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CHARACTERIZATION OF CHROMOPHORES IN ELECTRONICALLY  
EXCITED STATES BY MEANS OF COMPUTATIONAL METHODS

CHARAKTERIZACE CHROMOFÓRŮ V ELEKTRONICKY EXCITOVANÝCH  
STAVECH POMOCÍ VÝPOČETNÍCH METOD

Habilitační práce

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## **Poděkování**

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Jsem si vědoma, že bez dlouholeté podpory vedení Ústavu Organické chemie a biochemie, AV ČR bych tuto práci nemohla předložit.

## **Prohlášení**

Prohlašuji, že jsem tuto habilitační práci zpracovala samostatně a že jsem uvedla všechny použité informační zdroje a literaturu.

V Praze, dne 22.4.2024

## Obsah

1. Úvod
    - 1.1. Charakterizace excitovaných stavů
    - 1.2. Terminologie pro popis excitovaných stavů
    - 1.3. Excimery, exciplexy a excitony
    - 1.4. Neadiabatické procesy
    - 1.5. Výpočetní metodika
  2. Fotochemie nukleových kyselin
    - 2.1. Fotochemie monomerů
    - 2.2. Vliv okolních bází na absorpční spektra
    - 2.3. Vliv okolních bází na procesy v excitovaných stavech
  3. Excitované stavy v kvantových tečkách
    - 3.1. Charakterizace excitovaných stavů
    - 3.2. Efekt dopování heteroatomy na optické vlastnosti CD
    - 3.3. Interakce mezi povrchy polyaromatických uhlovodíků
  4. Shrnutí
- Reference
- Příloha: Publikace s autorčíným příspěvkem, titulní strana

## 1. Úvod

Habilitační práce předkládá výsledky studií, jejichž společným tématem je charakterizace excitovaných stavů. Modelování excitovaných stavů je ilustrováno na dvou typech systémů: (1) polyaromické uhlovodíky a odvozené systémy jako modely kvantových teček v základním a excitovaných stavech a (2) excitované stavy izolovaných systémů bází nukleových kyselin a jejich modifikací, a bází interagujících pomocí nekovalentních interakcí. Práce, na kterých se autorka podílela, jsou specificky zmíněny a zařazeny do kontextu tehdejších studií na dvou uvedených systémech. Jevy, na jejichž studiu se autorka nepodílela, jsou zmíněny pouze okrajově z důvodu zachování tematické celistvosti.

### *1.1. Charakterizace excitovaných stavů*

Určení elektronové struktury molekuly v základním stavu je prvním krokem procesu její charakterizace, včetně chemických vlastností a reaktivity. Charakter elektronové struktury převážně většiny molekul v základním stavu dovoluje popis systému pomocí jedné elektronové konfigurace (jedno-referenční popis). Vhodným ukazatelem je dostatečný energetický rozdíl mezi 'highest occupied molecular orbital' (HOMO) a 'lowest unoccupied molecular orbital' (LUMO). Pro některé systémy a procesy, např. existuje-li více energeticky blízkých elektronových konfigurací (systémy s d-orbitaly ve valenční slupce, radikály), reakce molekul zahrnující tvorbu nebo přerušení chemické vazby a procesy v excitovaných stavech, je tento popis nedostatečný. Tyto systémy se nazývají multi-referenční a pro jejich popis je třeba zahrnout další elektronové konfigurace.<sup>1</sup>

Popis procesů probíhajících v excitovaných stavech navíc zahrnuje situace, ve kterých selhává jedna ze základních aproximací kvantové chemie, Born-Oppernheimerova aproximace (BOA) separující elektronový a jaderný pohyb. Tato separace je klíčovým předpokladem řešení Schrödingerovy rovnice a spolu s Franck-Condonovým principem, který říká, že k elektronickému přechodu s největší pravděpodobností dojde beze změn polohy jader v molekule a jejím okolí, tvoří základ pro výpočet elektronických spekter.<sup>2-4</sup> Mimo Franck-Condonovu oblast může dojít k situaci, kdy se k sobě dva nebo více povrchů potenciální energie (PES) energeticky přiblíží, a tak se pohyby jader a elektronů propojí a BOA selhává. Tato situace otevírá témata fotoindukovaných procesů,<sup>5</sup> neadiabatické fotodynamiky<sup>6</sup> a neadiabatických interakcí.<sup>7,8</sup> Aparát pro popis neadiabatických procesů bude ilustrován v dalším textu.

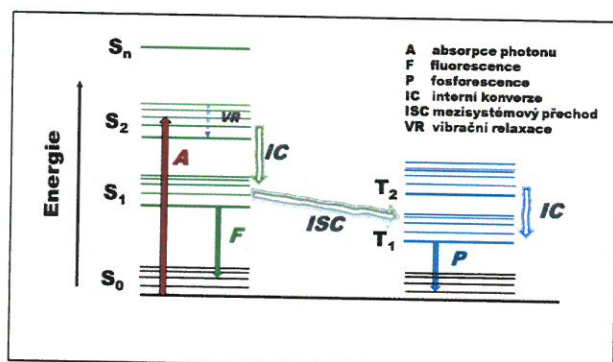
Molekulární procesy indukované UV/VIS zářením hrají významnou roli v mnoha vědeckých a technologických oborech. Pro experimentální i teoretické chemiky je odlišné chování molekul v

elektronicky excitovaných stavech ve srovnání se základním stavem skutečnou výzvou. Experimentální metody, jmenovitě konvenční stacionární spektroskopie<sup>9</sup> a její rozšíření na časově rozlišenou spektroskopii,<sup>10</sup> jsou schopny popsat systémy v excitovaných stavech s velkou přesností. Pro interpretaci těchto dat jsou vzhledem ke komplexnímu chování molekul v excitovaných stavech výpočetní metody nezbytným prostředkem.

### 1.2 Terminologie pro popis excitovaných stavů

Elektronové stavy molekuly a přechody mezi nimi shrnuje Jablonského diagram (Obr. 1). Absorpce elektromagnetického záření **A** UV nebo viditelných vlnových délek vede k vybuzení molekuly do nerovnovážného stavu, z kterého se spustí sekvence relaxačních procesů.

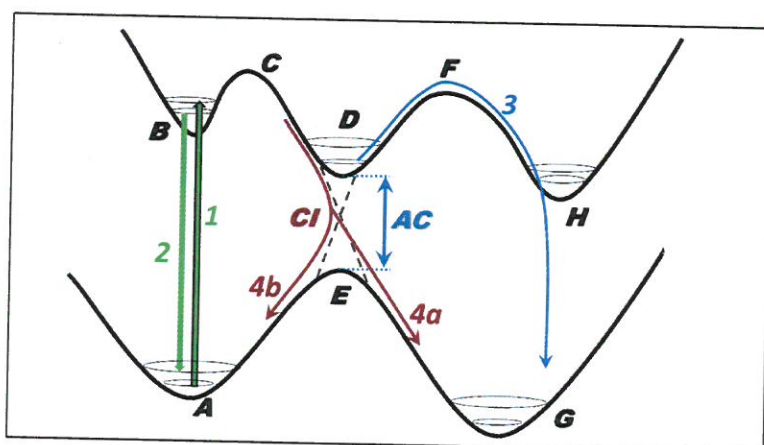
Tyto procesy mohou odpovídat vibrační relaxaci **VR** na jedné PES, neradiačnímu přechodu na jinou adiabatickou PES stavu se stejnou (vniřní konverze, **IC**) nebo rozdílnou spinovou multiplicitou (mezisystémový přechod, **ISC**), případně radiačním přechodům, tj. fluorescenci **F** nebo fosforescenci **P**. Relaxační procesy probíhají v časovém rozpětí v řádu několika desítek femtosekund, což odpovídá vibračním relaxacím, až milisekund, odpovídající fosforescenci, tedy v časovém rozpětí více než deseti řádů.



Obr. 1. Jablonského diagram definující procesy: absorpce fotonu (A), fluorescence (F), fosforescence (P), interní konverze (IC), mezisystémový přechod (ISC), vibrační relaxace (VR)

Po absorpci světla v bodě **A** systém přejde do bodu **B** ve Franckově-Condonově oblasti (proces 1 na Obr. 2). Po nalezení minima excitovaného stavu se systém může vrátit do základního stavu emisí světla (proces 2 na Obr. 2) nebo zůstat v excitovaném stavu. V druhém případě se systém vyvíjí ve třech následujících procesech: (i) setrvává po nějakou dobu v excitovaném stavu, kde může případně interagovat s jinými excitovanými stavy (viz níže), (ii) následně relaxuje do základního stavu a (iii) sleduje PES základního stavu do minima. Jak znázorňuje Obr. 2, při dostatečně velké počáteční energii (proces 1) systém překoná bariéru transičního stavu **C** a přejde k blízkosti minima excitovaného stavu, tj. do bodu **D**, z kterého může následovat dvě možné cesty. V prvním případě se povrchy v místě pseudo-křížení základní a excitované PES (body **D** a **E**) nekříží (tzv. avoided crossing, **AC** na Obr. 2), systém setrvá na excitované PES a přes transiční stav **F** přejde do dalšího

minima excitovaného stavu **H**, odkud relaxuje do základního stavu **G**, které není identické s **A**, tzn. vzniká tzv. fotoprodukt. Tento proces, znázorněný křivkou **3**, probíhá až do bodu **H** beze změny elektronického stavu a je označován za adiabatický. V druhém případě se křivky kříží v bodě zvaném kónická intersekcce (conical intersection, **CI** na Obr. 2), relaxační proces probíhá na dvou PES a odpovídající mechanismus je nediabatický. Jak vyplývá z obr. 2, systém může relaxovat jak do minima **G** za vzniku fotoproduktu (křivka **4a**), tak do původního minima **A** (křivka **4b**). V druhém případě mluvíme o fotostabilním systému. Na rozdíl od adiabatického mechanismu, kdy je systém v minimu **D** charakterizován konečnou dobou života, procesy probíhající v **CI** jsou mimořádně rychlé, v časovém intervalu molekulových vibrací, a systém tak není možné v tomto stavu experimentálně zachytit.



Obr. 2. Adiabatické a nediabatické procesy na křivkách potenciální energie dvou stavů.

### 1.3 Excimery, exciplexy a excitony

V případě absorpce UV záření systému s větším množstvím absorbujících skupin (chromoforů) může dojít k jejich interakci v excitovaných stavech. Tyto interakce jsou výsledkem dvou hlavních příspěvků:

- (1) Coulombické interakce, tzv. 'long-range', které lze popsat interakcí transitních dipólových momentů nebo vyšších multipólů a závisí na vzdálenosti chromoforů ( $R$ ) jako  $R^{-3}$
- (2) tzv. 'short-range' interakce, které jsou výsledkem překryvu orbitalů zapojených do excitace. Dosah těchto interakcí klesá exponenciálně s  $R$ .

Podle povahy excitace lze excitované komplexy popsat v termínech lokalizovaných a 'charge-transfer' (CT) přechodů:

$$\Psi_{exc} = c_1\psi(A^*B) + c_2\psi(AB^*) + c_3\psi(A^-B^+) + c_4\psi(A^+B^-), \quad (1)$$

kde  $\psi(A^*B)$  a  $\psi(AB^*)$  odpovídá excitacím lokalizovaným na A nebo B a  $\psi(A^-B^+)$  a  $\psi(A^+B^-)$  popisuje CT stavy, při kterých dochází k přenosu náboje.<sup>11</sup> Je-li A identické s B a  $c_1 = c_2$ ,  $c_3 = c_4$ , výsledný komplex se nazývá excimerem, v opačném případě se jedná o exciplex.

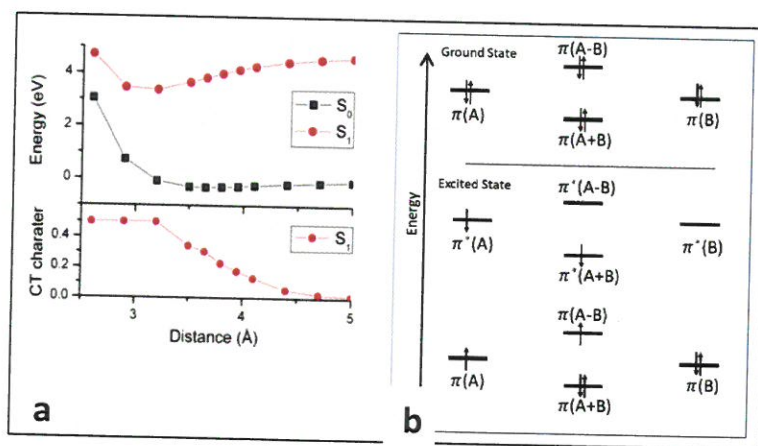
Povahu a rozsah interakce a tvorbu excimeru/exciplexu názorně ilustruje Obr. (3a) na příkladu dvou naphthalenů v symetrickém uspořádání.<sup>11</sup> V základním stavu tvoří dimer slabě vázaný komplex s mělkým minimem v  $R \cong 3.65 \text{ \AA}$  (získané výpočtem ADC(2) metodou) se stabilizační energií menší než 0.1 eV. V excitovaném stavu je stabilizační energie 1.32 eV s minimem v  $R \cong 3.2 \text{ \AA}$ . Povaha interakce v dimeru v excitovaném stavu je ilustrována příspěvků elektronových konfigurací s CT charakterem (Obr. 3a, spodní graf). V  $R > 5 \text{ \AA}$  jsou CT příspěvky nevýznamné a dimer má charakter tzv Frenkelova excitonu, který lze popsat Frenkelovou excitonovou teorií, podle které je celkový Hamiltonián systému vyjádřen výrazem

$$H = \sum_A H_A + \sum_A \sum_{B>A} V_{AB}, \quad (2)$$

kde  $H_A$  odpovídá izolovaným chromophorům a  $V_{AB}$  je interagující člen. Vlnovou funkci excitonu vyjadřuje tvar

$$\Psi_k = \sum_A c_{kA} \phi_A^i, \quad (3)$$

kde  $\phi_A^i$  je vlnová funkce excitovaného chromoforu A.



Obr. 3 a) Horní graf: Závislost stability naphthalene dimer v základním ( $S_0$ ) a excitovaném ( $S_1$ ) stavu na vzdálenosti, Spodní graf: příspěvek CT charakteru k celkové vlnové funkci.<sup>11</sup> b) Molekulární diagram tvorby excimeru



Je-li  $R < 5 \text{ \AA}$ , CT stavy přispívají ke stabilizaci systému (pro naphthalene dimer je příspěvek v minimu 50%), systém získává povahu excimeru a výsledná funkce je kombinací CT a Frenkelova excitonu.

Větší stabilizaci komplexu lze vysvětlit pomocí molekulárního diagramu orbitalů uplatňujících se při excitaci (např.  $\pi$  orbitalů) na Obr. 3b. Při tvorbě dimeru A-B jsou  $\pi$  a  $\pi^*$  orbitály výsledkem symetrické kombinace ( $\pi(A+B)$  a  $\pi^*(A+B)$ ) a antisymetrické kombinace ( $\pi(A-B)$  a  $\pi^*(A-B)$ ) orbitalů s energií nižší (vazebná, + kombinace) nebo vyšší (antivazebná, - kombinace) ve srovnání s orbitálními energiemi monomerů. V základním stavu jsou  $\pi(A+B)$  a  $\pi(A-B)$  plně obsazeny, a tedy dva elektrony obsadí stabilizující vazebný orbital a dva elektrony obsadí destabilizující antivazebný orbital. V celkové bilanci je příspěvek nulový a k stabilizaci přispívají pouze disperzní interakce. V excitovaném stavu obsadí tři elektrony stabilizující kombinaci a jeden elektron destabilizující kombinaci orbitalů, což vede k větší stabilitě oproti základnímu stavu.

#### 1.4 Neadiabatické procesy

Pro popis neadiabatických jevů vedoucích k silné interakci mezi různými elektronovými konfiguracemi zúčastněných stavů je třeba zavést formalismus, který zohledňuje propojení pohybů jader a elektronů. Následující text stručně popisuje příslušný formalismus.

Časově nezávislá Schrödingerova rovnice pro systému má tvar

$$H^T \Psi^T(r, R) = E^T \Psi^T(r, R), \quad (4)$$

kde  $R$  je souřadnice jader,  $r$  je souřadnice elektronů a Hamiltonian  $H^T$  je definován

$$H^T(r, R) = T^N(R) + H(r, R), \quad (5)$$

kde  $T^N(R)$  je operátor kinetické energie jader a  $H(r, R)$  je elektronový Hamiltonián, který obsahuje kinetickou energii elektronů a Coulombickou interakci mezi částicemi. Řešením Schrödingerovy rovnice pro elektrony

$$H \Psi_I(r, R) = E_I(R) \Psi_I(r, R), \quad (6)$$

jsou vlastní funkce  $H$ , které mohou být použity k rozvoji celkové vlnové funkce ve formě

$$\Psi^T(r, R) = \sum \psi_I(r, R) \chi_I(r, R), \quad (7)$$

kde  $\chi_I$  jsou rozvojové koeficienty. Použitím tohoto rozvoje lze po úpravách získat rovnici

$$[T^N + E_I] \chi_I - \sum \sum \frac{1}{2M_\alpha} (2f_{Ij}^\alpha \nabla_\alpha \chi_j + K_{Ij}^\alpha \chi_j) = E^T \chi_I, \quad (8)$$

kde  $\alpha$  a  $J$  se vztahují ke stupňům volnosti jader a elektronů, respektivě. Členy  $K_{IJ}$  a  $f_{IJ}$  popisují interakci elektronických stavů  $I$  a  $J$ , vyplývají z operátoru kinetické energie jader operujícího na elektronické vlnové funkce  $\psi_I(r, R)$  a mají tvar

$$f_{IJ}^\alpha(R) = \langle \psi_I(r, R) | \nabla_\alpha \psi_J(r, R) \rangle \quad (9)$$

$$K_{IJ}(R) = \sum \langle \psi_I(r, R) | \nabla_\alpha^2 \psi_J(r, R) \rangle, \quad (10)$$

kde  $\nabla_\alpha$  definuje gradient vzhledem k souřadnicím jader  $R$ .

Ve většině případů je systém popsán jedním elektronovým stavem, členy  $K_{IJ}$  a  $f_{IJ}$  jsou zanedbatelné a mohou tak být ignorovány, což vede k tzv. Born-Oppenheimerově aproximaci. V oblasti pseudo-křížení a křížení dvou PES musí být tyto členy zahrnuty. Vektor  $f_{IJ}$  je dominující, je měřítkem změny elektronické vlnové funkce vzhledem ke změnám souřadnic jader a je nepřímo úměrný rozdílu energií stavů  $I$  a  $J$ :

$$f_{IJ}(R) = \frac{\langle \psi_I | \nabla H | \psi_J \rangle}{E_I - E_J} \quad (11)$$

Je-li celková funkce  $\Psi^T$  rozvinuta v bázi vlastních funkcí celkového Hamiltoniánu  $H^T$ , je interakce mezi jednotlivými elektronovými stavy zahrnuta v členu pro kinetickou energii jader  $T^N$  a nediagonální členy matice  $H^T$  jsou rovné nule. Taková reprezentace se nazývá adiabatická. Alternativně lze použít diabatickou reprezentaci  $\phi_I$ , ve které je interakce elektronických stavů vyjádřena nenulovými hodnotami nediagonálních členů matice Hamiltoniánu, který má tvar

$$H^{diab} = \begin{bmatrix} H_{11} & H_{12} \\ H_{21} & H_{22} \end{bmatrix} \quad (12)$$

a vlastní funkce mají tvar

$$E_{1,2} = \bar{H} \pm \sqrt{\Delta H^2 + H_{12}^2}, \quad (13)$$

kde  $\bar{H} = (H_{11} + H_{22})/2$  a  $\Delta H = (H_{11} - H_{22})/2$ .

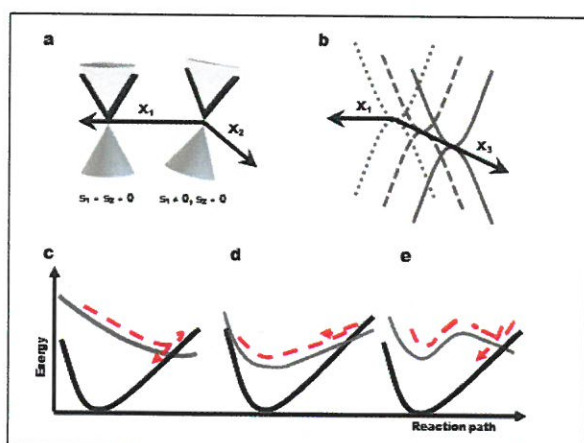
Pro degenerované hodnoty  $E_1$  a  $E_2$ , tj. křížení dvou PES, musí být splněny podmínky  $H_{11} - H_{22} = 0$  a  $H_{12} = 0$ . Tyto podmínky jsou splněny pro  $N^{\text{int}} - 2$  podprostorů, kde  $N^{\text{int}}$  je počet stupňů volnosti molekuly. Dvouatomová molekula má pouze jeden stupeň volnosti, a tak pro dva elektronické stavy o stejném spinu a symetrii tato podmínka nemůže být splněna, což vede k tzv. pravidlu o nekřížení stavů. Polyatomické molekuly splňují tuto podmínku pro  $N^{\text{int}} - 2$  stupňů volnosti. Po jistých úpravách lze matici (12) přepsat v souřadnicích kónické intersekcce:

$$H = (s_1x + s_2y)\mathbf{I} \begin{bmatrix} gx & hy \\ hy & -gx \end{bmatrix} \quad (14)$$

kde  $x$  a  $y$  odpovídají posunu ve směru vektoru gradientu energetického rozdílu  $\mathbf{g}$  o normě  $g$  a vektoru mísení (coupling)  $\mathbf{h}$  o normě  $h$ . Vektory  $\mathbf{g}$  a  $\mathbf{h}$  definují tzv. “branching space”, podprostory, ve kterých je degenerace porušena, ostatní orthogonální podprostory definují tzv. “crossing seam”.  $s_1$  a  $s_2$  jsou projekcemi  $(g_1+g_2)/2$  do roviny  $g$ - $h$ ,  $\mathbf{I}$  je 2x2 jednotková matice. Vlastní funkce Hamiltoniánu (14) mají tvar

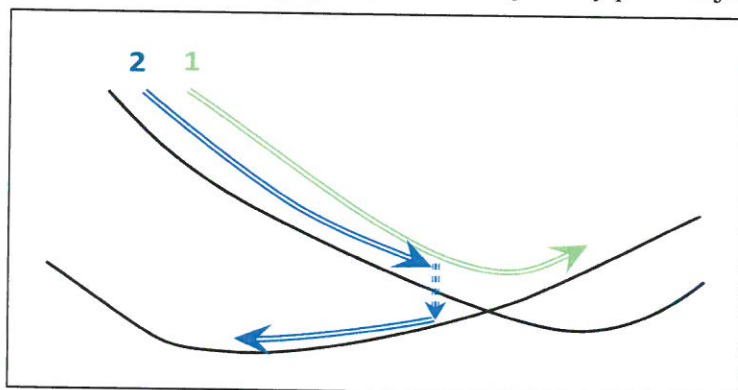
$$E_{1,2}(x, y) = s_1x + s_2y \pm \sqrt{(gx)^2 + (hy)^2} \quad (15)$$

Pomocí této rovnice je možné vykreslit PES obou stavů v okolí kónické intersekcce, která má tvar dvojitého kužele podél dvou souřadnic (Obr. 4). Parametry  $g$  a  $h$  udávají tvar **CI**, parametry  $s_1$ ,  $s_2$  udávají její naklonění, kde nulové hodnoty  $s_1$  a  $s_2$  odpovídají vertikální **CI**. Obr. 4a znázorňuje tvar PES podél souřadnic  $x_1$  a  $x_2$ , které definují “branching space”; Obr. 4b podél souřadnice  $x_3$ , která definuje “crossing seam”. Tvar **CI** a jejího okolí má zásadní vliv na její efektivitu pro neadiabatický přechod, a tedy jeho pravděpodobnost. Obr. 4c-4e znázorňují příklady **CI** o různé efektivitě v důsledku charakteru PES v její blízkosti. V případě, kdy PES excitovaného stavu monotónně klesá k bodu křížení (Obr. 4c), je **CI** nejefektivnější a pravděpodobnost přechodu z jedné PES na druhou je velká. V situaci, kdy PES excitovaného stavu sleduje stoupající trend směrem k **CI** (Obr. 4d), se otevírá trajektorie zpět do minima excitovaného stavu. Odpovídající **CI** je odkloněna z vertikálního směru, je méně efektivní a přechod je méně pravděpodobný. Podobně, existence bariéry na PES excitovaného stavu na cestě k **CI** (Obr. 4e) způsobuje nižší pravděpodobnost přechodu.



Obr. 4. Kónické intersekcce: (a) podél souřadnic  $x_1$  a  $x_2$ , které definují “branching space”, (b) podél souřadnice  $x_3$ , která definuje “crossing seam”. Typy reakčních cest excitovaného stavu vedoucích k **CI**: (a) monotónně klesající, (d) stoupající z minima excitovaného stavu, (e) překonávající bariéru.

Popis procesů, při kterých BOA selhává, vyžaduje použití metod schopných zohlednit rozdělení populace mezi několik elektronových stavů. Jednou z nejpoužívanějších metod pro simulaci neadiabatických jevů je tzv. 'trajectory surface-hopping' (TSH) metoda. V rámci této metody se struktura propaguje na jedné elektronické potenciální energetické ploše a při každém kroku stochastický algoritmus vyhodnocuje, ve kterém stavu bude trajektorie pokračovat (obr. 5). Tento postup se opakuje dokud není trajektorie dokončena. Vzhledem k stochastické povaze metody TSH je popis dynamiky plnohodnotný pouze v případě, že se uvažuje velký počet trajektorií.



Obr. 5. Dvě možné situace v TSH metodě: (1) – adiabatická trajektorie a (2) – neadiabatická trajektorie

### 1.5 Výpočetní metodika

Pro modelování systémů vyžadujících zahrnutí více elektronových konfigurací je třeba použít multi-konfigurační metodu (multi-configurational self-consistent field, MCSCF). V rámci této metody je celková vlnová funkce konstituována jako lineární kombinace Slaterových determinantů (SD), každý popisující jednu elektronovou konfiguraci. MCSCF tak kombinuje metodu konsistentního pole (Self Consistent Field, SCF) a konfigurační interakce (Configuration Interaction). Výsledná energie systému je získána optimalizací dvou typů rozvojových koeficientů: koeficientů  $c_{si}$  lineárního rozvoje atomových orbitalů  $\chi_s$  při optimalizaci molekulových orbitalů  $\varphi_i$  (rovnice (16), MO-LCAO) v SCF

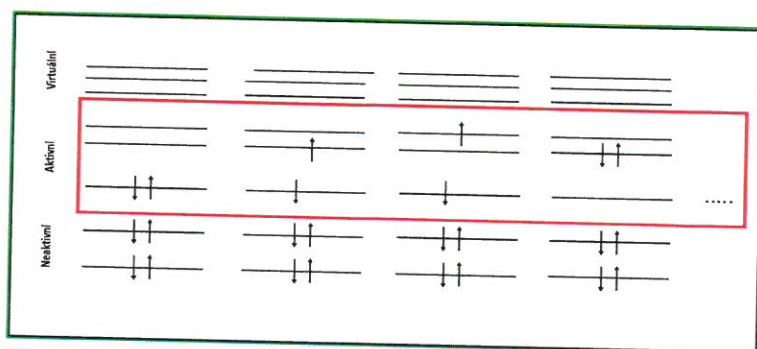
$$\varphi_i(1) = \sum_s^M c_{si} \chi_s(1), \quad (16)$$

(symbol (1) vyjadřuje, že každý atomový orbital, a tedy i molekulový orbital, závisí na souřadnicích jednoho elektronu (elektronu 1)) a koeficientů  $d_l$  lineárního rozvoje SD  $\Phi_l$  v rozvoji konfigurační interakce

$$\Psi = \sum_I d_I \Phi_I. \quad (17)$$

Speciálním případem MCSCF metody je tzv. Complete Active Space Self-Consistent Field (CASSCF) metoda, ve které je omezen výběr elektronových konfigurací rozdělením orbitálního prostoru na neaktivní, aktivní a virtuální (Obr. 6). V neaktivním prostoru jsou orbitály obsazeny dvěma elektrony (DOCC), v aktivních orbitalech tvořících tzv. aktivní prostor (AS) je možná jakákoliv elektronová konfigurace. Vybraný prostor je charakterizován počtem elektronů a orbitalů jako CAS( $n$  elektronů,  $m$  orbitalů). Orbitály v DOCC i AS jsou předmětem optimalizace.

Pro diskusi MCSCF metody je vhodné zavést pojem elektronové korelace, tedy vzájemného vlivu pohybu elektronů. Elektronová korelace bývá někdy rozdělována na dynamickou a statickou (nedynamickou). Dynamická korelace je důsledkem pohybu elektronů, nedynamická je spojena s nutností zahrnout do vlnové funkce systému více elektronových konfigurací, jak je popsáno v MCSCF metodě. Ostrá hranice mezi statickou a dynamickou korelací neexistuje.



Obr.6. Schéma molekulových orbitalů pro tvorbu aktivního prostoru.

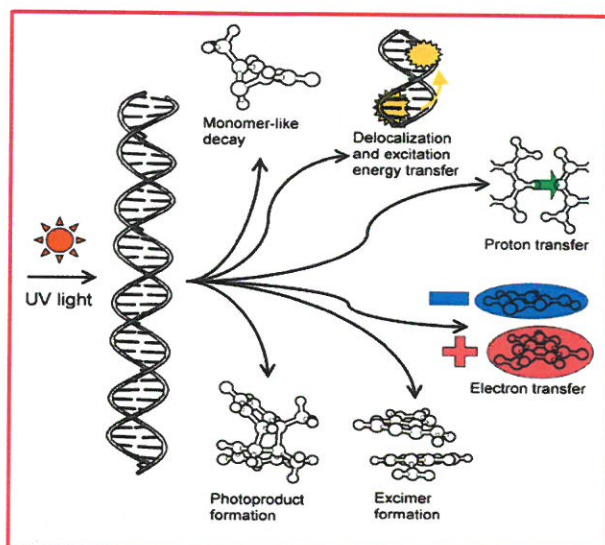
## 2. Fotochemie nukleových kyselin

Komplexita chování nukleových kyselin (DNA i RNA) v excitovaném stavu představuje velkou výzvu pro experimentální i výpočetní chemiky studující tyto procesy. Všechny báze nukleových kyselin (nukleobáze) absorbují UV záření a jejich excitace může iniciovat fotoreakce, které ve svém důsledku mohou vést k poškození genetického kódu, tzv. mutacím.<sup>12</sup> Hlavním úkolem je pochopení relaxačních procesů následujících po elektronické excitaci. Vedle zářivé relaxace existuje v nukleových kyselinách několik nezářivých relaxačních procesů (obr. 7):

- 1) Lokalizace elektronické excitace na jedné bázi, po níž následuje relaxace v rámci tohoto monomeru
- 2) Delokalizace excitovaných stavů vedoucí k tvorbě excitonů
- 3) Přenos energie elektronické excitace.
- 4) Tvorba silně interagujícího komplexu - excimeru

- 5) Přenos elektronů nebo protonů mezi dvěma řetězci nebo v rámci jednoho řetězce
- 6) Vznik fotochemického produktu

Reakce (6) neodpovídá situaci, kdy je zachována dříve zmíněná fotostabilita nukleových kyselin, a nebude proto dále diskutován.

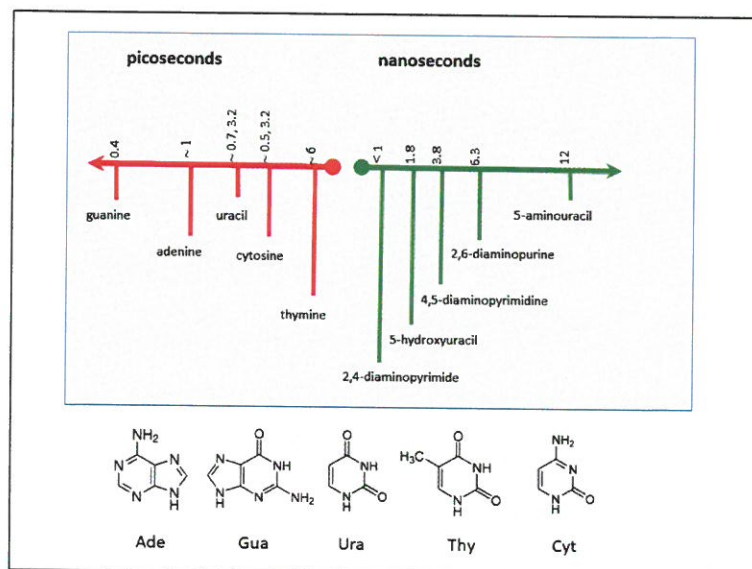


Obr. 7. Relaxační mechanismy v molekulách nukleových kyselin

### 2.1 Fotochemie monomerů

Báze nukleových kyselin patří do skupiny aromatických heterocyklických sloučenin odvozených z pyrimidinu (cytosin, thymin a uracil) a purinu (adenin a guanin), chromoforů silně absorbujících v UV oblasti. Díky přítomnosti heteroatomů jsou nízko položené excitované stavy výsledkem  $\pi \rightarrow \pi^*$  a  $n \rightarrow \pi^*$  přechodů, při vyšších energiích je možné dosáhnout také přechody do  $\sigma^*$  orbitalů. Variabilita excitovaných stavů mezi heterocykly založené na pyrimidinu a purinu je především ve stabilitě, tedy době života excitovaných stavů. Z obr. 8 je patrné, že společnou vlastností všech diskutovaných sloučenin vyskytujících se v nukleových kyselinách (nukleobázích) je mimořádně krátká doba života excitovaných stavů, v řádu picosekund.<sup>13</sup> Zároveň je známo, že tyto báze mají velmi nízký kvantový výtěžek luminescence. Nukleobáze jsou tak schopné relaxovat do základního stavu nezářivými nediabatickými mechanismy dříve než dojde k jejich chemické přeměně<sup>14-16</sup>, Click or tap here to enter text.tj. vykazují fotostabilitu danou pravděpodobně evolučním tlakem v raném biotickém věku. Velmi malé modifikace v chemické struktuře způsobí změnu doby života excitovaných stavů v několika řádech. Např. substitucí vodíku v uracilu nebo methylové skupiny v thyminu v pozici C5 za OH nebo NH<sub>2</sub> skupinu se doba života prodlouží z pikosekund na

nanosekundy (obr. 8).<sup>17</sup> Podobně i změna konformace má vliv na dobu života. Příkladem je cytosin a guanin, v kterých mají keto- formy kratší dobu života než enol formy.<sup>13,18</sup>



Obr. 8. Doby života excitovaných stavů bází nukleových kyselin a jejich modifikací

Výpočetní metody společně s ultrarychlou časově rozlišovanou spektroskopií<sup>12,19,20</sup> přispěly k značnému pokroku v pochopení neadiabatických procesů. Nicméně, interpretace experimentálních časově rozlišených spekter se v některých případech rozcházel v závislosti na použité metodice.<sup>21</sup>

*Autorka se podílela na komplexní studii všech nukleobází. Pomocí simulací neadiabatickou dynamikou na základě multi-referenčního popisu (CASSCF metodou) se podařilo získat sjednocující popis fotostability všech nukleobází<sup>22-27</sup>, a na základě srovnání s jejich modifikacemi<sup>17,28</sup> objasnit mechanismus jejich relaxace.* Výsledky těchto studií a jejich srovnání s experimentálními daty<sup>21</sup> jsou prezentovány na obr. 9 a v Tab 1. a pro jednotlivé báze ukazují následující:

*Adenin:* Experimenty provedené při absorpční energii ( $E_{\text{abs}}$ ) 250-270 nm ukazují, že doba života se zkracuje s rostoucí energií. Tento jev byl reprodukován pomocí simulací neadiabatickou dynamikou, ve které byly doby života 0.77 ps a 0.63 ps pro  $E_{\text{abs}}$  267nm a 250 nm. V obou případech relaxují všechny trajektorie s jedinou časovou konstantou, což ukazuje na homogenní relaxační mechanismus, A1, vedoucí k CI mezi  $\pi\pi^*$  excitovaným a základním stavem (closed-shell, cs), charakterizované distorzí C<sub>2</sub> atomu z roviny purinového cyklu.

*Guanin:* z experimentálních studií prováděných s  $E_{\text{abs}}$  267 nm byly v závislosti na způsobu fitování navrženy doby života 0.36 ps a 0.8 ps. Dynamické simulace reprodukovaly relaxaci s kratší dobou života. Také v tomto případě systém relaxuje s jedním relaxačním mechanismem G1 vedoucím k CI( $\pi\pi^*/\text{cs}$ ), ve které je C<sub>2</sub> atom distortován z roviny purinového cyklu.

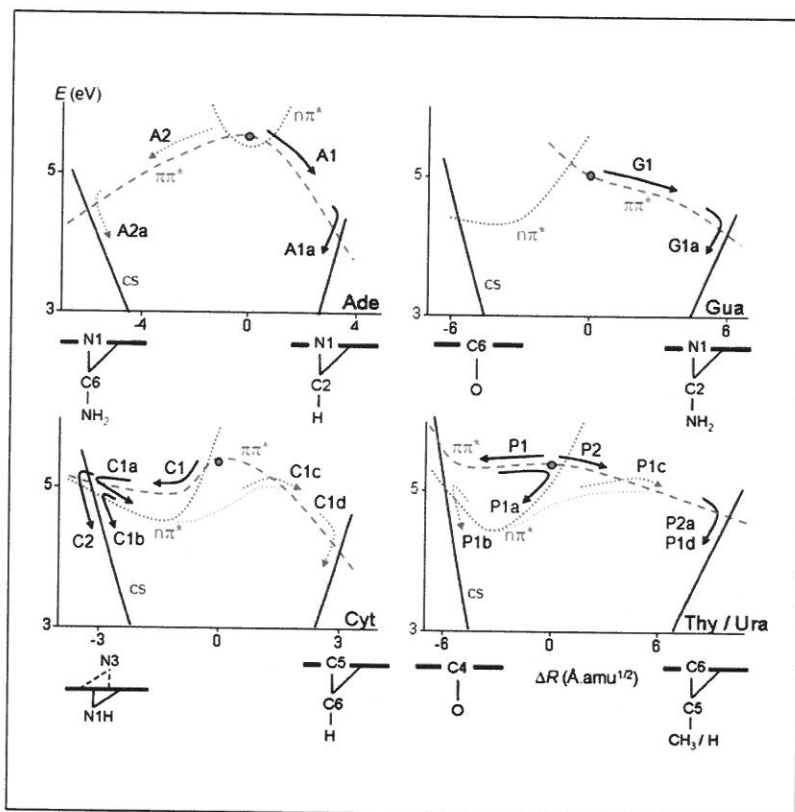
*Cytosin*: Experimentální studie provedené s  $E_{\text{abs}} = 250$  nm poskytly doby života 0.8 ps a 3.2 ps. Experimenty s  $E_{\text{abs}} = 267$  nm, v blízkosti centra prvního absorpčního bandu, nebyly publikovány. Podle dynamických simulací 74% populace relaxuje s časovou konstantou 0.53 ps, 13% s časovou konstantou 3.08 ps, což je v dobré shodě s výsledky experimentu. Simulace navíc ukázaly, že 13% populovaných stavů relaxuje v ultra-krátké době během 10 fs. Relaxační schéma cytosinu je mnohem komplexnější ve srovnání s guaninem a adeninem. Relaxační mechanismy sledované pro trajektorie s časovou konstantou 0.53 ps následují v 87% trajektorii s přechodem z  $\pi\pi^*$  do  $\pi\pi^*$  stavu přes quasilanární CI (C1a). Z této CI se mechanismy dělí na dráhu C1b (58%), vedoucí ke quasilanární CI na intersekcí  $\pi\pi^*/\text{cs}$ , a na dráhu C1c vedoucí na C1d a dále k CI ( $\pi\pi^*/\text{cs}$ ), charakterizované distorzí C6 atomu z roviny pyrimidinového kruhu, a následné relaxaci do základního stavu. Oba mechanismy probíhají na stejné časové škále. Minoritní frakce (13%) relaxuje během 10 fs na dráze C2.

*Thymin*: Představuje nukleobázi s nejdelší dobou života, v řádu několika ps, jak bylo pozorováno v experimentech s  $E_{\text{abs}} = 267$  nm a 250 nm. Simulace provedené při těchto excitačních energiích ukázaly časové konstanty 2.3 ps v případě vyšších energií a 1.7 ps při nižších energiích. Po relaxaci do  $S_1(\pi\pi^*)$  (P1a) stavu jsou dvě možnosti: 1. relaxace pomocí CI na intersekcí  $\pi\pi^*/\text{cs}$  stavů (P1b), charakterizovanou distorzí O atomu z roviny pyrimidinového kruhu, 2. přechod do  $\pi\pi^*$  stavu a dále k CI na intersekcí  $\pi\pi^*/\text{cs}$  stavů (P1c a P1d), charakterizovanou distorzí C5 atomu z roviny pyrimidinového kruhu. Vzhledem k limitované době simulace 3 ps nelze jednoznačně predikovat, který mechanismus převažuje, i když nižší energie CI( $\pi\pi^*/\text{cs}$ ) naznačuje, že P1c mechanismus převládá.

*Uracil*: Experimenty byly provedeny s  $E_{\text{abs}} = 250$  nm, a to na dvou časových škálách, 0.5 ps a 2.4 ps. Při nižších energiích ( $E_{\text{abs}} = 267$  nm) byla pozorována pouze relaxace s delší dobou života. Většina trajektorií (77%) simulována při nižší počáteční energii zrelaxovala během 1.8 ps z  $S_2$  do  $S_1$  stavu, další relaxace přesahovala čas simulace, tj. 2 ps. 23% trajektorií se zrelaxovalo do základního stavu během 0.74 ps. Podobné výsledky dávaly i simulace při vyšší  $E_{\text{abs}}$ , i když v kratším časovém úseku a s vyšší populací plně relaxujících trajektorií. Analýza trajektorií navrhla tři mechanismy relaxace. Většina populace (60%) relaxuje obdobným mechanismem jako thymin do  $\pi\pi^*$  minima (P1a). Stejně jako u thyminu, následný relaxační mechanismus nebylo možno určit v rámci doby simulace. Druhý mechanismus, objevující se pro 31% trajektorií, který přispívá k deaktivaci s časovou konstantou 0.74 ps/0.66 ps, používá trajektorii podél  $\pi\pi^*$  stavu (P2) k



CI( $\pi\pi^*/cs$ ) charakterizované twistem podél C5C6 vazby. K deaktivaci probíhající na stejné časové škále přispívá i třetí mechanismus spojený s otevřením pyrimidinového kruhu mezi N3 a C4 atomy. Srovnání relaxačních mechanismů purinových a pyrimidinových bází ukazuje, že adenin a guanin



Obr. 9. Relaxační mechanismy bází nukleových kyselin získané pomocí neadiabické dynamické studie.<sup>21</sup>

představují fotodeaktivací mechanismus následující diabaticky  $\pi\pi^*$  stav přímo a bez bariéry k CI na intersekcí se základním stavem. K dosažení struktur lokalizovaných v CI je potřeba určitá koordinace v pohybech všech atomů, které ve svém výsledku vedou k distorzi purinového systému z roviny, vedoucí k poněkud odlišným časům pro adenin a guanin. Fotodynamika pyrimidinových bází je mnohem komplexnější a jejich doba života je větší, více než 0.5 ps. Důvodem je poměrně mělké minimum na PES  $\pi\pi^*$  stavu způsobené existencí několika stavů  $\pi\pi^*$  různého charakteru a delokalizovanými  $\pi$  orbitaly. Existence mělkého minima zpomaluje 'nalezení' přímé  $\pi\pi^*$  cesty. Detaily tvaru PES a energetických bariér pak určují další rysy dynamiky.

Dynamické simulace jednotlivých bází ukazují, že neexistuje žádná jedinečná fotodynamická deaktivace, nicméně lze definovat specifické vlastnosti pro určité relaxační cesty. Za zmínku stojí především déle trávající (několik ps) výskyt pyrimidinů, zvláště thyminu v  $S_2$  minimu, tzv 'trapping'. Podobné zpomalovací mechanismy mohou s určitou pravděpodobností hrát roli v dalších fotoindukovaných procesech, např. tvorby pyrimidinových dimerů,<sup>29</sup> při které se předpokládá

uplatnění tripletových excitovaných stavů. Singlet-tripletní přechod probíhá na větší časové škále a nemůže konkurovat relaxacím probíhajícím v rámci několika pikosekund, nicméně zpomalení podobná výše uvedenému trappingu by existenci tripletních stavů mohl umožnit. Podpůrným argumentem pro roli tripletních stavů v relaxaci nukleobází jsou studie naznačující, že v přítomnosti vody je aktivován  $n\pi^*$  stav a z něj následně triplet.<sup>30</sup>

Bližší pochopení dynamického chování nukleobází přináší také studie alternativních systémů, které ukazují na vliv jednotlivých strukturálních parametrů, především distorzí z roviny kruhu nebo prodlužování vazeb, a jejich změn na pravděpodobnosti jednotlivých relaxačních cest, a tedy efektivity příslušných CI. Takové situace byly studovány např. na 2,4-diaminopyrimidinu a 2,6-diaminopurinu.<sup>31</sup> Podobně, velmi malé modifikace v chemické struktuře způsobí změnu doby života excitovaných stavů v několika řádech. Např. substitucí vodíku v uracilu nebo methylové skupiny v thyminu v pozici C5 za OH nebo NH<sub>2</sub> skupinu se doba života prodlouží z pikosekund na nanosekundy (obr. 8).<sup>17</sup>

Tab.1 Časové konstanty relaxace nukleobází

	$\lambda$ (nm)	$\tau_1$ (ps)		$\tau_2$ (ps)		$\tau_3$ (ps)		Reference <sup>a</sup>
		Exp	Exp	Simulated	Exp	Simulated		
Ade	250	< 0.05	0.75	0.63				1, 14
	267	0.10	1.0	0.77				2, 15-19
	277		-2					14,17
Gua	267	0.15	0.36 (0.86) <sup>b</sup>	0.28				2, 16
Thy	250	<0.05	0.49		6.4	>2.3		1
	267	0.11			5.1	>1.7		2,16
Ura	250	<0.05	0.53	0.65	2.4	>1.5		1
	267	0.13		0.74	1.1	>1.8		2,16
Cyt	250	<0.05	0.82		3.2			1
	267	0.16		0.53	1.86 (3.2) <sup>b</sup>	3.08		2,16

<sup>a</sup>reference v článku PNAS, <sup>b</sup>rozdílné hodnoty jsou dány rozdílnou fitovací procedurou v experimentu

## 2.2 Vliv okolních bází na absorpční spektra

Znalost povahy excitovaného stavu v nukleových kyselinách ve smyslu lokalizace excitace na monomeru nebo delokalizace je důležitá pro charakterizaci následného vývoje excitovaného systému. V molekulách nukleových kyselin jsou možné dva typy delokalizace, a to tvorba excitonu nebo CT excitovaných stavů.<sup>11</sup> *Autorka se podílela na studiích kratších fragmentů<sup>32,33</sup> ve vzájemných orientacích odpovídajících rozdílným konformacím DNA*, které ukázaly, že vliv okolních bází, a tedy i strukturálních fluktuací, na excitační energie monomeru je minimální, podstatně ale ovlivňuje interakci transičních dipóvých momentů, tedy interakci mezi monomery.

***Předmětem studií, na kterých se autorka podílela, bylo i ověření spolehlivosti výpočetních metod pro interpretaci absorpčních spekter.*** Bylo ukázáno, že v souladu s obecně známým faktem, že TDDFT metoda přeceňuje stabilitu CT stavů, metody založené na vlnové funkci ukazují, že tyto stavy jsou méně energeticky stabilní než lokalizované stavy. Navíc mají CT stavy ve srovnání s lokalizovanými stavy menší transiční dipolový moment (sílu oscilátoru), a tak menší příspěvek v absorpčních spektrech nukleových kyselin.<sup>34</sup>

### *2.3 Vliv okolních bází na procesy v excitovaných stavech*

Efekty, které ovlivňují fotoaktivitu nukleobází, mohou být rozděleny na vnější (vliv okolního prostředí) a vnitřní (elektronické interakce mezi bázemi). Vzhledem k velikosti systému (až několik desítek bází) a požadavkům na výpočetní metodiku (multireferenční popis) je modelování těchto efektů výpočetně velmi náročné. Obvyklým přístupem je molekulová dynamika zahrnující neadiabatické efekty kombinující QM metody pro popis bází účastnících se excitace s MM metodami modelující okolní subsystém, tzv QM/MM metoda. ***Autorka se podílela na pracech studujících jak vliv 'stackových' interakcí a vodíkových vazeb v rámci jednoho vlákna,<sup>26</sup> tak vodíkových vazeb mezi bázemi sousedních vláken.<sup>27</sup>*** Výsledky ukázaly, že zejména vodíkové vazby ovlivňují dosažitelnost CI purinových bází, charakterizovaných distorzí kruhu z roviny. Zatímco v prvním případě mají vodíkové vazby tendenci urychlovat relaxační proces podpořením této distorze, v druhém případě omezuje distorzi, a tím zabraňuje přístup k CI, jak bylo ukázáno na studiu dynamiky Gua v Cyt-Gua páru bází v DNA. Vodíková vazba neovlivnila dobu života excitovaného Cyt, relaxace Gua se ale významně zpomalila. Příčinou je existence vodíkové vazby mezi kyslíkovým atomem C=O vazby cytosinu a vodíkovým atomem NH<sub>2</sub> skupiny guaninu, která omezuje dostupnost CI zodpovědnou za ultrakrátkou dobu života, kratší než 0.5 ps.

Tvorba excimerů/excipleků, neutrálních i CT, je příkladem relaxace excitovaných stav v důsledku elektronových interakcí mezi bázemi, tedy vlivem vnitřních efektů. ***Autorka přispěla ke studiu těchto komplexů analýzou elektronické interakce z hlediska Coulombických a orbitálních překryvových členů.<sup>32,33</sup>*** V těchto pracech bylo ukázáno, že elektronická interakce je výsledkem především Coulombických interakcí, zatímco orbitální překryv při vzdálenosti odpovídající vzdálenosti ve šroubovici DNA významně nepřispívá. Ukázalo se rovněž, že Coulombické interakce mezi bázemi se vzájemnou orientací odpovídající konformaci A-DNA je větší než v případě která odpovídá konformaci B-DNA a že změny v excitačních energiích způsobených nesymetrickými účinky jsou o řád větší než ty, které jsou způsobeny "čistými" elektronickými efekty. Vliv dimerizace včetně možné tvorby excimerů byly studovány také na xanthinových derivátech, systémech podobných nukleobázím.<sup>35</sup>

Výčet relaxačních mechanismů doplňuje tzv. 'proton coupled electron transfer' (PCET) probíhající na jedné z vodíkových vazeb WC báze<sup>36-41</sup> předpokládající velmi rychlý proces deaktivace. Tento proces nebyl předmětem studia autorky, a proto není podrobněji diskutován.

### 3. Excitované stavy v kvantových tečkách

Od svého objevení se uhlíkové tečky (carbon dots, CD),<sup>42</sup> nanočástice na bázi uhlíku s alespoň jedním rozměrem menším než 10 nm, rozšířily do velké třídy nekovových fotoluminescenčních nanomateriálů. Jejich nespornou výhodou je nízká nebo nulová toxicita, biokompatibilita, dispergovatelnost v široké škále rozpouštědel, biologická odbouratelnost, nízké náklady, snadná a škálovatelná syntéza, předurčují CD k použití v mnoha oblastech životního prostředí. Díky své biokompatibilitě jsou CD ideálními kandidáty pro biosenzory v zobrazovacích metodách<sup>43,44</sup>, jako nosiče pro podávání léčiv a protinádorovou terapii.<sup>45,46</sup>

CD jsou malé fragmenty grafenu s konečnou velikostí, čímž vykazují na rozdíl od grafenu nenulový energetický gap mezi obsazenými a neobsazenými orbitály, a tak jedinečné optické a elektronické vlastnosti, včetně výrazné fotoluminescence (PL). Díky variabilní velikosti, možnosti dopování heteroatomy, funkcionalizaci povrchu a přítomnosti různých typů defektů je možné dosáhnout PL od UV přes viditelnou až po blízkou infračervenou oblast a vykazovat různé PL charakteristiky, jako např. kvantový výtěžek, posuny v PL spektrech, atd. Za PL jsou zodpovědné tzv. jádra obsahující konjugované  $\pi$ -domény, celkové PL chování ale ovlivňuje také chemická struktura na okrajích CD. Elektronový gap s rostoucí velikostí klesá, což vede k červenému posunu emisního pásu. Jako u jiných systémů, i v CD je možná tvorba excimerů a excitonů nebo nezářivé relaxace a přechody na tripletní stavy, což výrazně ovlivňuje PL vlastnosti. Výčet faktorů, které mohou ovlivnit tyto vlastnosti doplňují další parametry, jako detaily struktury, chemické modifikace, způsob syntézy, reakční podmínky, atd. Popis PL vlastností a interpretace experimentálních dat tak představuje mimořádně komplikovaný problém, který je stále intenzivně studován.<sup>35</sup>

CD jsou běžně modelovány pomocí polyaromatických uhlovodíků (polyaromatic hydrocarbons, PAH), tzv. polyaceny, circumaceny a periaceny. V závislosti na struktuře a velikosti mohou mít vlnové funkce PAH v základním stavu uzavřenou nebo otevřenou (radikálovou) slupku a/nebo podíl excitovaných elektronických konfigurací. Charakteristickým rysem je přítomnost několika rezonančních struktur, jejichž příspěvky k vlnové funkci základního stavu rostou s velikostí systémů. Charakter vlnové funkce lze poměrně snadno a spolehlivě určit pomocí Clarova pravidla,<sup>35,47-49</sup> podle kterého má struktura s více aromatickými sextety (tzv. 'Clar' sextet) vyšší aromatickou stabilizační energii, díky čemuž je možné, že radikálová struktura je stabilnější než struktura s uzavřenou slupkou (Kekulé struktura). Obr 10 ilustruje existenci dvou rezonančních

struktur heptacenu (Obr 10a): Kekulé strukturu s uzavřenou slupkou a jedním 'Clar' sexteterm a radikálovou strukturu s dvěma 'Clar' sextety; a 3-2-3-periacenu (Obr. 10b): Kekulé strukturu s uzavřenou slupkou a dvěma 'Clar' sextety a biradikálovou strukturou se čtyřmi 'Clar' sextety.

*Autorka se podílela na dílčích problémech spojených s charakterizací těchto systémů v základním i excitovaném stavu, které jsou uvedeny v následujícím textu.*

### 3.1 Charakterizace excitovaných stavů

Radikálová povaha základního stavu spolu s příspěvky elektronových konfigurací s dvojnásobnou excitací pro větší systémy, jakož i multi-referenční povaha excitovaných stavů komplikuje výběr metodiky pro popis optických vlastností CD. Nejpoužívanější metodou pro tyto účely je jednoznačně metoda založená na tzv. 'time-dependent density functional theory' (TDDFT). Uvedenou metodu, respektivě vhodnost jednotlivých funkcionalů, je nutné díky výše uvedeným problémům týkajících se základního stavu některých CD pečlivě testovat srovnáním s experimentem a metodami schopných správně popsat základní stav a následně poskytnout správné pořadí (ve smyslu energie) excitovaných stavů. *Autorka se podílela na několika pracech, které se touto problematikou zabývají. Studie zahrnovaly charakterizaci excitovaných stavů serie polyacenů, naphthalenu ( $n=2$ ,  $n$  je počet benzenových kruhů) až heptacenu ( $n=7$ ),<sup>50</sup> ve které byly detailně studovány statické i dynamické příspěvky korelační energie v jednotlivých excitovaných stavech.* Pro větší systémy bylo ukázáno, že vzhledem ke kvazi-degenerovanému charakteru vlnových funkcí nemusí být běžně používané DFT metody vhodnou volbou a pro výpočty velkých PAH by se měly přednostně používat multireferenční přístupy.

Přesnost jednotlivých funkcionalů pro výpočty excitovaných stavů, především určení správného pořadí stavů jak lokalizovaných, tak CT stavů, bylo testováno na modelových systémech: pyrene, circum-1-pyrene, coronene, circum-1-coronene, and circum-2-coronene (obr 10c).<sup>51</sup> Absorpční spektra byly počítány multireferenčními metodami, 'density functional theory/multireference configuration interaction' (DFT/MRCI), 'spectroscopy oriented configuration interaction' (SORCI) a 'contracted-nelectron valence state perturbation theory to second order' (NEVPT2), a jednoreferenčními metodami 'scaled opposite-spin-algebraic diagrammatic construction to second-order' (SOS-ADC(2)), TD-B3LYP a TD-Coulomb-attenuating method-B3LYP (TD-CAM-B3LYP). Srovnáním s dostupnými experimentálními daty bylo ukázáno, že nejkonzistentnější metodou je DFT/MRCI a je doporučena pro menší systémy. Pro větší systémy, kde vzhledem k výpočetní náročnosti nelze použít multireferenční metody, jsou ze sledovaných metod nejvhodnější SOS-ADC(2) a CAM-B3LYP.

### 3.2 Efekt dopování heteroatomů na optické vlastnosti CD

Implementace heteroatomů do CD struktur je jedna z možností jak významně ovlivnit optoelektronické vlastnosti.<sup>52</sup> Např. v závislosti na koncentraci a poloze heteroatomu lze získat rozdílné typy posunů, tj. ke kratším i delším vlnovým délkám. Přítomnost heteroatomů také zvyšuje kvantový výtěžek fluorescence. Dopování N atomem, kterého lze dosáhnout během syntézy s použitím prekurzorů obsahujících atomy dusíku, je nejčastější.<sup>53-55</sup>

***Autorka se podílela na studii zabývající se vlivu pozice N(H)-dopací pyrenu (1, obr. 10d). Byly zkoumány tzv. 'N-grafitic' (2-4, obr. 10d), 'N-grafitic-edge' (5 a 6, obr. 10d) a NH-pyridinové (7 – 9, obr. 10d) dopace na stabilitu elektronových struktur v základním stavu a vliv přítomnosti N atomů na optická spektra CD.***<sup>53,56</sup> Analýza provedená na základě multireferenčních výpočtů a kvalitativně (teorie valenční vazby, Clarovo pravidlo, HOMA index pro hodnocení aromaticity atd.) umožnila racionalizovat relativní stability různých izomerů a také příslušné posuny v UV spektrech. V 'N-grafitic' a 'N-grafitic-edge' struktuře se počet  $\pi$  elektron ve srovnání s pyrenem zvyšuje o dva. V případě pyridinových pozic zůstává počet  $\pi$  elektronů nezměněn. V případě grafitického dopování se elektronický charakter řídí vzájemnou polohou atomů dusíku. Jsou-li N atomy lokalizovány na stejném typu alterujících atomů (uhlíkové atomy neoznačené nebo označené '\*' v obr. 10d, struktura 1), má struktura silný biradikálový charakter a není možné zkonstruovat žádnou Kekulé strukturu. Přítomnost Clarova sextetu tento efekt ještě zesiluje. Bylo také zjištěno, že struktury s větší vzdáleností N atomů (ve smyslu počtu uhlíkových atomů mezi atomy dusíku) jsou stabilnější. Pyridinové substituce neovlivňují optická spektra ve srovnání s pyrenem na rozdíl od grafitických pozic, u kterých charakter absorpčních spekter závisí na výše diskutované (ne)radikálovém charakteru základního stavu. V případě radikálové povahy, tj. N atomy jsou v pozicích na stejném typu alterujících atomů a odděluje je tedy lichý počet atomů C, jsou  $S_0 \rightarrow S_1$  excitační energie v rozmezí od 0.5 – 1 eV, tj. spektrum je výrazně posunuto k nižším energiím (červený posun) do infračervené oblasti, zatímco grafitický N na okrajích GQDs může vyvolat buď červený nebo modrý posun.

***Autorka se rovněž podílela na studiu vlivu přítomnosti F atomů na optická spektra CD.***<sup>57</sup> V uvedené práci byl zkoumán vliv fluoru vázaného kovalentní vazbou, tj. na okrajích CD struktur, a vliv aniontem fluoru na absorpční a emisní spektra pyrenu a circum-1-pyrenu. Především byla zkoumána možnost dosažení červeného posunu ve spektrech a zvýšení intenzity pásů. Studie byla provedena pomocí DFT/MRCI, ADC(2) a DFT/CAM-B3LYP metod. Jako v případě N-dopování, také vliv F-dopování závisí na pozici heteroatomu. Je-li F lokalizován nad aromatickým kruhem za vzniku záporně nabitého komplexu, dochází k výraznému posunu k nižším energiím v absorpčním i emisním spektru. Tyto změny jsou způsobeny vznikem nového excitovaného stavu s velkým

podílem CT. Tento jev nebyl dříve pro komplexy tohoto typu zvažován. Kovalentní substituce nemají zásadní vliv na povahu absorpčních spekter. Navíc, optimalizace do minima excitovaného stavu ukazuje, že F má tendenci disociovat.

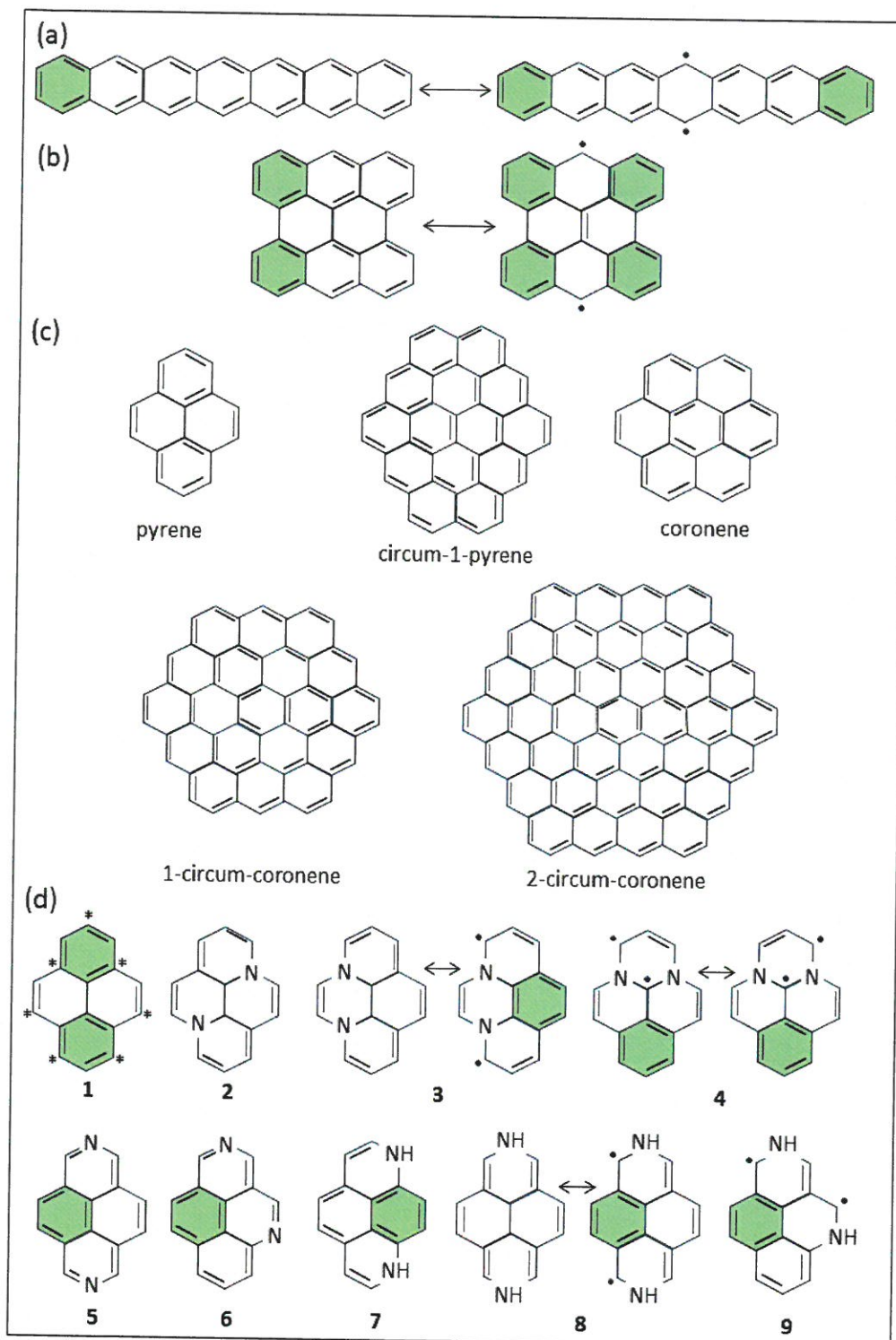
### 3.3 Interakce mezi PAH povrchy

V důsledku agregace mohou jednotlivé vrstvy PAH vzájemně interagovat pomocí tzv. 'stacking' interakcí. Tyto interakce vyvolávají změny optických vlastností podobných těm pozorovaných v excimerech/exciplech Jejich možné využití bylo naznačeno v několika studiích.<sup>58,59</sup>

K požadavkům na výpočetní metodu popsaných dříve pro monomery se přidávají další požadavky, především popis CT mezi jednotlivými monomery, popis elektronové korelaci na delší vzdálenost a ve srovnání se základním stavem větší polarizabilita systému ovlivňující interakce v excitovaném stavu. Pro popis excimerů/exciplech je vhodnost standardní TD-DFT problematická, a to ze dvou důvodů: díky nesprávnému popisu disperzních interakcí nezbytných k získání přesných vazebných energií v základním a excitovaném stavu pro slabě vázané van der Waalovy komplexy a díky nadhodnocené stabilitě stavů s CT charakterem.<sup>60</sup>

K modelování těchto asociátů se obvykle používají menší aromatické systémy (benzen, pyrene, perylene, atd.)<sup>60,61</sup> *Autorka se podílela na studiu absorpčních a emisních spekter rozšířených modelů založených na pyrenu a coronenu,<sup>60,62</sup> ve které byla rovněž testována spolehlivost TD-DFT (konkrétně TD-CAM-B3LYP) a ADC(2) metod na základě srovnání s DFT/MRCI výpočty a experimentálními daty.* Studie ukázala, že obě jednoreferenční metody poskytují uspokojivou shodu s 'benchmark' daty pro nižší-energetické excitované stavy, TD-CAM-B3LYP má ale tendenci podhodnocovat CT stavy, a tak nesprávně predikovat pořadí excitovaných stavů. I přes tyto problémy poskytuje TD-CAM-B3LYP rozumnou predikci změny v emisních spekter se zvětšujícím se počtem aromatických kruhů ve srovnání s DFT/MRCI (obr. 11). Rovněž extrapolace vedla k podobným hodnotám emisní energie pro nekonečně velké systémy ve srovnání s ADC(2) metodou, tj. 1.63 eV (760 nm, ADC(2)) a 1.77 eV (700 nm, TD-CAM\_B3LYP), tedy hodnotám blížíci se IČ oblasti.

Charakter emisních spekter, především otázka, zda je určován jedním chromophorem nebo agregáty více chromophorních skupin je jednou z hlavních otázek týkajících se optických vlastností CD. Zajímavým fenoménem pozorovaným experimentálně je výrazný rozdíl mezi absorpčními a emisními energiemi,<sup>63,64</sup> tzv. Stoke shift.



Obr. 10. Příklady rezonančních struktur polyaromatických uhlovodíků. Clar sextety jsou vyznačeny zeleně

*Z Tab 2 publikované v práci,<sup>62</sup> na níž se autorka podílela, je patrné, že hodnota Stokesova posunu klesá s rostoucí velikostí systému a dosahuje limitní hodnoty pro nekonečně velké systémy. To ukazuje, že experimentálně pozorované Stokesovy posuny nemohou být vysvětleny v rámci*

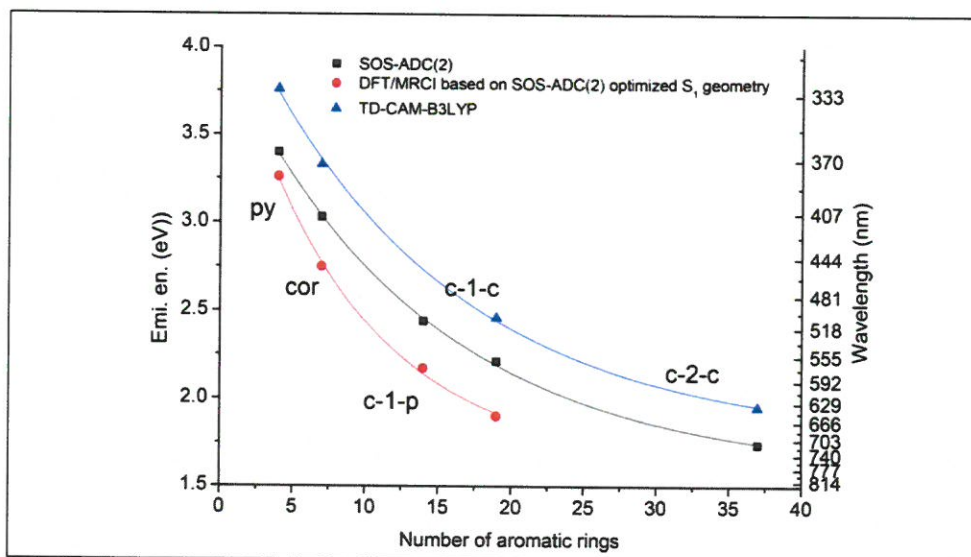


monomerů a jsou s největší pravděpodobností důsledkem agregátů, jak potvrzují také hodnoty získané pro dimery.

Tab. 2 Stokes shift (eV) pyrenu, circum-1-pyrenu, coronenu, circum-1-coronenu, circum-2-coronenu počítaný ADC(2) a TD-CAM-B3LYP metodami a def2-TZVP bází<sup>a</sup>

	TD-CAM- B3LYP	ADC(2)	
	Monomer	Monomer	Dimer
pyrene	0.27	0.31	0.96
circum-1-pyrene	-- <sup>b</sup>	0.21	0.72
coronene	0.27	0.29	0.46
circum-1-coronene	0.16	0.17	0.53
circum-2-coronene	0.12	0.15	

<sup>a</sup>pro circum-2-coronene byla použita SVP báze, <sup>b</sup>nebylo vyhodnocováno v důsledku špatného řazení excitovaných stavů



Obr. 11 Závislost emisních energií na velikosti systémů polyaromatických uhlovdíků (počtu fúzaných benzenových kruhů)

#### 4. Shrnutí

Předkládaná práce shrnuje výsledky studií dvou biologicky relevantních systémů

1. Studium procesů v elektronicky excitovaných stavech bází nukleových kyselin, a to jak izolovaných, tak v prostředí tvořeném sousedními nukleobázemi. Tyto studie jsou významné k pochopení mimořádné fotostability nukleobází uplatňující se velmi pravděpodobně v raných fázích vzniku života, tak pochopení principů, kterými se nukleové báze chrání před svým poškozením vlivem, např. slunečního záření. Autorka se podílela na komplexní studii relaxačních mechanismů nukleobází jak v izolovaném stavu, tak interagujících s ostatními bázemi v rámci struktury nukleových kyselin. K pochopení mechanismu přispěly rovněž

autorčiny práce popisující chování excitovaných stavů systémů, které představují modifikace nukleobází. V studiích byly použity metody založené na přesném *ab initio*, převážně multi-referenčním popisu systémů. Výsledky lze shrnout do následujících bodů:

- i) Izolované nukleobáze relaxují rozdílnými mechanismy pro ty založené na purinu a pyrimidinu. Purinové báze mají kratší dobu života a při relaxaci do základního stavu sledují jeden  $\pi\pi^*$  stav vedoucí bezbariérově ke kónické intersekcí se základním stavem. Relaxace pyrimidinových bází je komplexnější v důsledku mělkého minima, která zpomaluje 'nalezení' přímé  $\pi\pi^*$  cesty. Další rozdíly dynamiky nukleobází jsou pak způsobeny detaily tvaru PES a energetických bariér.
  - ii) Modifikace nukleobází mají vliv na strukturní charakter kónických intersekcí, např. distorzí z roviny kruhu neb prodlužování vazeb, a a tak na dosažitelnost, a tedy efektivitu příslušných CI.
  - iii) Okolní báze ovlivňují relaxační mechanismus jednak možnou tvorbou eximerů, ve kterých se uplatňují tzv. 'stacking' interakce v rámci jednoho vlákna nukleových kyselin, tak interakcemi v rámci dvou vláken, které mohou omezit strukturální změny nutné k dosažení příslušné kónické intersekcí.
2. Charakterizace uhlíkových teček, které vzhledem ke své netoxicitě a biologické kompatibilitě a zároveň výrazným fotoluminescenčním vlastnostem představují materiály s vysokým potenciálem využití v medicíně. Na druhou stranu, v důsledku komplexního charakteru těchto systémů představuje racionalizace jejich fotoluminescenčních vlastností stále velkou výzvu. Výpočetní modelování je komplikováno také velikostí uvedených systémů, a tak problémy s použitím přesných výpočetních metod. Studie provedené na souboru polyaromatických uhlovodíků a jejich modifikovaných systémech ukázaly:
- i) Bezpečnou metodou pro spolehlivou interpretaci absorpčních spekter je multireferenční metoda schopná správně popsat jak 'statickou', tak 'dynamickou' korelaci. Použití TD-DFT metod vyžaduje pečlivé testování funkcionalů.
  - ii) Vliv přítomnosti dopantů na povahu absorpčních spekter, konkrétně výrazných posunů k nižším absorpčním energiím byl vysvětlen na základě analýzy elektronové povahy základního stavu u N-dopovaných PAH a na přítomnosti nově vzniklých 'charge-transfer excitovaných stavů u F-dopovaných PAH.
  - iii) Bylo ukázáno, že experimentálně pozorovaný posun mezi absorpčními a emisními pásy je s největší pravděpodobností způsoben přítomností agregátů vzniklých mezi chromoforními skupinami.

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## Příloha

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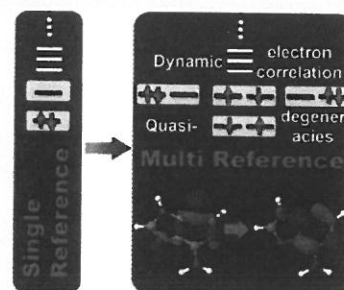
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## Multireference Approaches for Excited States of Molecules

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**ABSTRACT:** Understanding the properties of electronically excited states is a challenging task that becomes increasingly important for numerous applications in chemistry, molecular physics, molecular biology, and materials science. A substantial impact is exerted by the fascinating progress in time-resolved spectroscopy, which leads to a strongly growing demand for theoretical methods to describe the characteristic features of excited states accurately. Whereas for electronic ground state problems of stable molecules the quantum chemical methodology is now so well developed that informed nonexperts can use it efficiently, the situation is entirely different concerning the investigation of excited states. This review is devoted to a specific class of approaches, usually denoted as multireference (MR) methods, the generality of which is needed for solving many spectroscopic or photodynamical problems. However, the understanding and proper application of these MR methods is often found to be difficult due to their complexity and their computational cost. The purpose of this review is to provide an overview of the most important facts about the different theoretical approaches available and to present by means of a collection of characteristic examples useful information, which can guide the reader in performing their own applications.



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## CONTENTS

1. Introduction	7294	2.3.1. Single Reference Coupled Cluster Methods To Describe Excited States	7301
2. Theory and Methods	7295	2.3.2. Multireference Coupled Cluster Methods	7302
2.1. Configuration Interaction	7295	2.4. Multireference Perturbation Theory	7305
2.1.1. Basic Concepts	7295	2.5. MRCI with Semiempirical Hamiltonians	7307
2.1.2. Single- and Multireference Spaces	7295	2.6. Multireference Density Functional Theory	7308
2.1.3. Size-Extensivity Problem	7297	2.6.1. Nondynamical Electron Correlation in DFT	7308
2.1.4. Contracted CI	7297	2.6.2. Semiempirical MRCI with DFT	7308
2.1.5. Calculation of Excited States	7298	2.6.3. Hybrid Wave Function and DFT	7310
2.1.6. Individual Selection Schemes	7298	2.6.4. Multiconfigurational DFT	7310
2.1.7. Local MRCI Approaches	7298	2.6.5. Ensemble DFT	7310
2.1.8. Natural Orbitals for Use in MRCI	7298	2.7. Emerging Algorithms: DMRG and FCIQMC	7311
2.2. Multiconfiguration Self-Consistent Field Method	7299		
2.2.1. Multiple MCSCF Solutions and Symmetry Breaking	7299		
2.3. Coupled Cluster Methods	7300		

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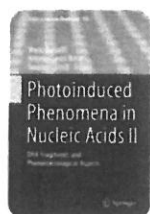
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Photoinduced Phenomena in Nucleic Acids

# Electronic Excitation Processes in Single-Strand and Double-Strand DNA: A Computational Approach

| Chapter I First Online : 01 January 2014

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## Photoinduced Phenomena in Nucleic Acids II

Felix Plasser, Adélia J.A. Aquino, Hans Lischka & Dana Nachtigallová

- **Part of the book series:** Topics in Current Chemistry ((TOPCURRCHEM, volume 356))

## Abstract

Absorption of UV light by nucleic acids can lead to damaging photoreactions, which may ultimately lead to mutations of the genetic code. The complexity of the photodynamical behavior of nucleobases in the DNA double-helix provides a great challenge to both experimental and computational chemists studying these processes. Starting from the initially excited states, the main question regards to understanding of the subsequent relaxation processes, which can either utilize monomer-like deactivation pathways or lead to excitonic or charge transfer species with new relaxation dynamics. After a review of photophysical processes in single nucleobases we outline the theoretical background relevant for interacting chromophores and assess a large variety of computational approaches relevant for understanding of the nature and dynamics of excited states of DNA. The discussion continues with the analyses of calculations on excitonic and charge transfer states followed by the presentation of the dynamics of excited-state processes in DNA. The review is concluded by topics of proton transfer in DNA and photochemical dimer formation of nucleobases.

Review


# Excited state dynamics of DNA bases

Karl Kleinermanns , Dana Nachtigallová & Mattanjah S. de Vries

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## Abstract

Biochemical reactions are subject to the particular environmental conditions of planet earth, including solar irradiation. How DNA responds to radiation is relevant to human health because radiation damage can affect genetic propagation and lead to cancer and is also important for our understanding of how life on earth developed. A reductionist approach to unravelling the detailed photochemistry seeks to establish intrinsic properties of individual DNA building blocks, followed by extrapolation to larger systems, to incorporate interactions between the building blocks and the role of the biomolecular environment. Advances in both experimental and computational techniques have lead to increasingly detailed insights in the excited state dynamics of DNA bases in isolation as well as the role of the solvent and intermolecular interactions. This review seeks to summarise current findings and understanding.

## The effect of C5 substitution on the photochemistry of uracil

Dana Nachtigallová,<sup>\*a</sup> Hans Lischka,<sup>\*ab</sup> Jaroslaw J. Szymczak,<sup>b</sup> Mario Barbatti,<sup>b</sup> Pavel Hobza,<sup>a</sup> Zsolt Gengeliczki,<sup>c</sup> Gustavo Pino,<sup>d</sup> Michael P. Callahan<sup>e</sup> and Mattanjah S. de Vries<sup>\*f</sup>

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A combined experimental and theoretical study on the excited-state behavior of the uracil analogues, 5-OH-Ura and 5-NH<sub>2</sub>-Ura is reported. Two-photon ionization and IR/UV double-resonant spectra show that there is only one tautomer present for each with an excited state lifetime of 1.8 ns for 5-OH-Ura and 12.0 ns for 5-NH<sub>2</sub>-Ura as determined from pump-probe experiments. The nature of the excited states of both species is investigated by means of multi-reference *ab initio* methods. Vertical excitation energies, excited state minima, minima on the crossing seam and reaction paths towards them are determined. Sizeable barriers on these paths are found that provide an explanation for the lifetimes of several nanoseconds observed in the experiment.

### Introduction

Gas-phase and computational studies of RNA and DNA bases have revealed unique photophysical properties that are sensitive to subtle structural differences. Nucleobase excited-state lifetimes depend on isomerism,<sup>1,2</sup> substitution,<sup>3,4</sup> tautomer form<sup>5,6</sup> and non-covalent interactions, such as hydrogen bonding with complementary bases<sup>7,8</sup> or with solvent molecules<sup>9</sup> and stacking<sup>10</sup> bonding with other nucleobases. Theoretical understanding describes the excited-state dynamics in terms of rapid internal conversion mediated by conical intersections.<sup>8,11–14</sup> This pathway offers a mechanism to diffuse electronic excitation by converting it to ground-state internal energy which can be safely transferred to the environment. One intriguing implication is the possibility that this mechanism could have affected prebiotic chemistry on an early Earth, thus influencing the eventual make-up of the biomolecular building blocks that form life as we know it today.

Conical intersections that couple potential surfaces require geometrical changes in the excited states. For example for 9-methylguanine the internal conversion pathway seems to involve a strongly bent amino group at position 2.<sup>15–17</sup> The pathway for cytosine appears to involve C=C twisting<sup>18,19</sup> with a low barrier for keto-cytosine (producing a short vibronic spectrum<sup>20</sup>) and a considerable energy gap to the S<sub>0</sub> state for enol-cytosine, producing an extensive vibronic

spectrum.<sup>20</sup> Therefore it should be no surprise that the potential energy landscapes that form conical intersections are affected by different substitutions. Excited-state dynamics should be especially sensitive to substitutions in the positions of the molecular frame that include coordinates along which the intersections are formed.

Among several works dealing with the effect of substitution in photoexcited nucleobases<sup>4,14,21–28</sup> we would like to draw attention to the recent work by Mburu and Matsika.<sup>4</sup> These authors have explored the effect of substitutions on purines and found that substitution at the C2 position decreases the energy of the first  $\pi\pi^*$  state considerably whereas substitution at the C6 position has a much smaller effect; the carbonyl group has in general a stronger effect than the amino group;  $n\pi^*$  states for all substituted purines are blue-shifted compared to purine. To understand the effect on pyrimidine bases they compared with 2-pyrimidinone (the ring structure of cytosine without the amino group) and found that the energy of the first bright excited state correlates strongly with the nature, position, and orientation of the substituent.

We have also recently explored the effect of amino substitutions in pyrimidines by studying 2,4-diaminopyrimidine by means of multi-reference *ab initio* methods.<sup>29</sup> The calculations of stationary points in the ground and excited states, minima on the S<sub>0</sub>/S<sub>1</sub> crossing seam and connecting reaction paths show that several paths with negligible barriers exist, allowing ultrafast radiationless deactivation if excited at energies slightly higher than the band origin. Even though the NH<sub>2</sub> group attached to the C2 position blocks the path to one of the lowest energy conical intersections, namely the one associated with ring puckering at C2, there are still other alternatives for almost barrierless access to conical intersections associated with deformations at the C5, C6, and N1 positions.

Substituted uracil compounds can be viewed as structural analogues of 2,4-diaminopyrimidine in which the C2 and C4 positions are substituted by oxygen rather than amino groups (see Fig. 1). We now report the effect of substitution at the C5

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# Relaxation mechanisms of UV-photoexcited DNA and RNA nucleobases

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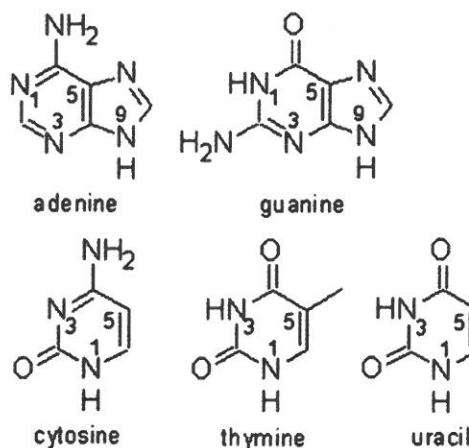
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A comprehensive effort in photodynamical *ab initio* simulations of the ultrafast deactivation pathways for all five nucleobases adenine, guanine, cytosine, thymine, and uracil is reported. These simulations are based on a complete nonadiabatic surface-hopping approach using extended multiconfigurational wave functions. Even though all five nucleobases share the basic internal conversion mechanisms, the calculations show a distinct grouping into purine and pyrimidine bases as concerns the complexity of the photodynamics. The purine bases adenine and guanine represent the most simple photodeactivation mechanism with the dynamics leading along a diabatic  $\pi\pi^*$  path directly and without barrier to the conical intersection seam with the ground state. In the case of the pyrimidine bases, the dynamics starts off in much flatter regions of the  $\pi\pi^*$  energy surface due to coupling of several states. This fact prohibits a clear formation of a single reaction path. Thus, the photodynamics of the pyrimidine bases is much richer and includes also  $n\pi^*$  states with varying importance, depending on the actual nucleobase considered. Trapping in local minima may occur and, therefore, the deactivation time to the ground state is also much longer in these cases. Implications of these findings are discussed (i) for identifying structural possibilities where singlet/triplet transitions can occur because of sufficient retention time during the singlet dynamics and (ii) concerning the flexibility of finding other deactivation pathways in substituted pyrimidines serving as candidates for alternative nucleobases.

photodynamical simulation | photostability | ultrafast photodeactivation | nonadiabatic interactions | *ab initio* multireference methods

Owing to the importance of mutagenic and carcinogenic effects caused by UV radiation on DNA, the UV-induced photochemistry and photophysics of individual bases (1–7), base pairs (8–10), and nuclei acid strands (11–13) have been intensively studied. In each of these levels, it has been found that the genetic code may count on molecular mechanisms to get rid of the energy excess minimizing deleterious effects. Isolated nucleobases (Scheme 1), for instance, are all photostable by returning to the ground state in an ultrafast time scale of a few picoseconds (1, 2, 14–20), which reduces the probability of undergoing photochemical transformations induced by reactive excited states. The ultrafast deactivation of the five natural nucleobases contrasts with the long lifetime of analogous bases not found in DNA and RNA (2, 21), suggesting that it could have been a source of evolutionary pressure in early biotic ages.

Ultrafast deactivation occurs through internal conversion mechanisms, where the molecule transfers the photoenergy stored in the electronic system to nuclear vibrational degrees of freedom. This means, the deactivation takes place without photoemission in contrast to what occurs for fluorescent and phosphorescent species. The energy transfer proceeds more intensely near nuclear geometries with state degeneracies also known as conical intersections, which motivates the search for mechanisms explaining how the molecular geometry is deformed from the ground state (or Franck–Condon) structure initially photoexcited



Scheme 1.

into the strongly distorted geometries where conical intersections usually occur.

Even though the nucleobases share the property of photostability, their structures, electronic properties, and spectra show considerable variation. Therefore, the question arises whether there is a common pattern in the deactivation dynamics of the nucleobases or does one encounter only a collection of processes that finally lead to the same result of photostability. Advanced techniques of theoretical chemistry are very well suited for performing such comparative studies even though the severe difficulties have to be recognized where several electronic states and their nonadiabatic interactions have to be considered at the same time. Together with ultrafast time-resolved spectroscopy (1, 2, 22), these methods have provided a high degree of detail in the characterization of internal conversion processes, suggesting ring puckering as one major structural theme for the photodeactivation of nucleobases (3, 23–26). Based on the available knowledge of computed structures and connecting reaction paths, a unifying description of the photostability of all nucleobases has been proposed (5, 23, 27). However, considering the fact that up to now for all nucleobases about 20 energetically accessible minimum energy conical intersection structures, which still describe the intersection seam only in a rudimentary way, are known, the goal of providing a theoretical explanation for the photodynamics based on description of reaction paths alone seems to be almost

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Photodynamics Simulations of Thymine: Relaxation into the First Excited Singlet State<sup>†</sup>Jaroslaw J. Szymczak,<sup>\*,‡</sup> Mario Barbatti,<sup>\*,‡</sup> Jason T. Soo Hoo,<sup>§</sup> Jaclyn A. Adkins,<sup>||</sup> Theresa L. Windus,<sup>‡</sup> Dana Nachtigallová,<sup>#</sup> and Hans Lischka<sup>\*,‡,#</sup>

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Ab initio nonadiabatic dynamics simulations are reported for thymine with focus on the  $S_2 \rightarrow S_1$  deactivation using the state-averaged CASSCF method. Supporting calculations have been performed on vertical excitations,  $S_1$  and  $S_2$  minima, and minima on the crossing seam using the MS-CASPT2, RI-CC2, MR-CIS, and MR-CISD methods. The photodynamical process starts with a fast ( $< 100$  fs) planar relaxation from the  $S_2 \pi\pi^*$  state into the  $\pi\sigma\pi^*$  minimum of the  $S_2$  state. The calculations demonstrate that two  $\pi$ -excited states (denoted  $\pi\pi^*$  and  $\pi\sigma\pi^*$ ) are actually involved in this stage. The time in reaching the  $S_2/S_1$  intersections, through which thymine can deactivate to  $S_1$ , is delayed by both the change in character between the states as well as the flatness of the  $S_2$  surface. This deactivation occurs in an average time of 2.6 ps at the lowest-energy region of the crossing seam. After that, thymine relaxes to the  $n\pi^*$  minimum of the  $S_1$  state, where it remains until the transfer to the ground state takes place. The present dynamics simulations show that not only the  $\pi\sigma\pi^*$   $S_2$  trapping but also the trapping in the  $n\pi^*$   $S_1$  minimum contribute to the elongation of the excited-state lifetime of thymine.

## 1. Introduction

Upon UV excitation, all five naturally occurring nucleobases return to the ground state on an ultrafast time scale ranging from half a picosecond to a few picoseconds.<sup>1–5</sup> The ultrafast decay minimizes the time that the molecule remains in reactive excited states, which could induce photochemical damage. This enhanced photostability might have been one factor favoring the selection of these bases in early biotic ages, to the detriment of other similar molecules with long-lived excited states.

In general, ultrafast decay depends on the existence of reaction pathways connecting the Franck–Condon region to the seam of conical intersections between the excited and ground states where radiationless processes can occur. For this reason, a great deal of theoretical work has been dedicated to the characterization of reaction paths and conical intersections not only for the five nucleobases<sup>6–18</sup> but also for their isomers,<sup>11,19</sup> substituted species,<sup>12,20,21</sup> and base models.<sup>22,23</sup> Further progress has been achieved by means of dynamics simulations, which attempt to describe the excited-state time evolution and the most frequently accessed reaction pathways. However, balancing the computational costs of dynamics simulations lasting for several picoseconds while still maintaining a proper description of multiple electronic excited states and their nonadiabatic couplings still constitutes a major challenge. Despite the difficulties, semiempirical,<sup>24–26</sup> density functional,<sup>27</sup> and ab initio<sup>28–31</sup> nonadiabatic dynamics simulations have been reported recently for nucleobases.

Among the five nucleobases, thymine has the longest lifetime.<sup>1,2</sup> Femtosecond-resolved pump–probe resonant ionization experiments pumped at 267 nm have revealed two exponential decay components, 6.4 and 100 ps, with the longer one assigned to the triplet-state population.<sup>2</sup> Another set of mass-selected femtosecond-resolved pump–probe resonant ionization experiments also pumped at 267 nm and identified a two-step mechanism with time constants 105 fs and 5.12 ps.<sup>1</sup> In the time-resolved photoelectron spectroscopy experiments reported in ref 4, three time constants were obtained,  $< 50$ , 490 and 6.4 ps (pump energy at 250 nm). Even though there is no full agreement about the details of the deactivation process, a time constant in the range of 5 to 6 ps for deactivation to the ground state clearly emerges from all of these experimental results. This time constant is, in addition, much larger than those measured for the other nucleobases (adenine: 1.1 ps, cytosine: 1.86 ps, guanine: 0.36 ps, uracil: 1.05 ps<sup>1</sup>).

On the basis of the reaction paths connecting the Franck–Condon region to the  $S_1/S_0$  conical intersections, Perun et al.<sup>13</sup> have proposed that the relatively long lifetime of thymine could be explained by a trapping of the molecule in the dark  $S_1 n\pi^*$  state after fast deactivation from the  $S_2 \pi\pi^*$  state. Nevertheless, multiple spawning dynamics simulations performed by Hudock et al.<sup>29</sup> at the complete active space self-consistent field (CASSCF) level have found a surprisingly small  $S_2 \rightarrow S_1$  deactivation yield in the first half picosecond. These authors then proposed that the reason for the long lifetime is the trapping of thymine in a  $S_2$  minimum right after the photoexcitation. On the basis of the analysis of the reaction paths connecting the minimum in the  $S_1$  state to the  $S_1/S_0$  conical intersections, Zechmann and Barbatti<sup>14</sup> have discussed how the low efficiency of those paths should be an additional factor adding to the  $S_2$  trapping to delay the deactivation to the ground state. Recently

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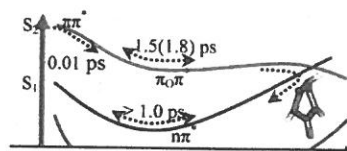


## Nonadiabatic Dynamics of Uracil: Population Split among Different Decay Mechanisms

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Supporting Information

**ABSTRACT:** Nonadiabatic dynamics simulations performed at the state-averaged CASSCF method are reported for uracil. Supporting calculations on stationary points and minima on the crossing seams have been performed at the MR-CISD and CASPT2 levels. The dominant mechanism is characterized by relaxation into the  $S_2$  minimum of  $\pi\pi^*$  character followed by the relaxation to the  $S_1$  minimum of  $n\pi^*$  character. This mechanism contributes to the slower relaxation with a decay constant larger than 1.5 ps, in good agreement with the long time constants experimentally observed. A minor fraction of trajectories decay to the ground state with a time constant of about 0.7 ps, which should be compared to the experimentally observed short constant. The major part of trajectories decaying with this time constant follows the  $\pi\pi^*$  channel and hops to the ground state via an ethylenic conical intersection. A contribution of the relaxation proceeding via a ring-opening conical intersection was also observed. The existence of these two latter channels together with a reduced long time constant is responsible for a significantly shorter lifetime of uracil compared to that of thymine.



## 1. INTRODUCTION

Upon UV excitation all five naturally occurring nucleobases relax into the ground state through internal conversion mechanisms on time scales not exceeding a few picoseconds,<sup>1–6</sup> in contrast to some of their analogues not found in nucleic acids.<sup>1,7</sup> This ultrafast decay serves as a protection mechanism, which prevents photochemical damage. This remarkable photostability of the nucleobases has been investigated intensively by means of experimental and theoretical approaches.

Ultrafast pump–probe techniques<sup>1–4,6–8</sup> have been applied in order to obtain detailed information on the specific time scales of the deactivation processes. For the purine bases, adenine and guanine,<sup>1,2,4,6,8,9</sup> experimental data indicate a single homogeneous mechanism with a faster decay compared to pyrimidine bases. The experimental studies performed on the pyrimidine bases, thymine, uracil, and cytosine, reveal more complicated photodynamics with generally longer lifetimes; the longest one is for thymine.<sup>1,2,4</sup> For uracil, using gas phase time-resolved photoelectron spectra obtained with a pump wavelength of 250 nm, three time constants were determined: a decay of less than 0.05 ps, followed by intermediary (0.53 ps) and longer (2.4 ps) decay constants.<sup>4</sup> Using the same initial energy, mass-selected femtosecond-resolved pump–probe ionization measurements were fitted with two exponentials, giving short and long decay times of 0.13 and 1.05 ps, respectively.<sup>1</sup> Kang et al. used a monoexponential fit in the analysis of femtosecond-resolved pump–probe resonant ionization spectra with 267 nm pump wavelength. They reported an exponential decay of 2.4 ps.<sup>2</sup> All these experimental results agree on a faster decay than that of thymine,<sup>1,2,4</sup> for which

lifetimes of 6.4 and 5.1 ps have been obtained. A dark state with a lifetime of several nanoseconds found in time-resolved photoelectron spectra of uracil and thymine was explained by trapping in the  $n\pi^*$  state.<sup>10,11</sup>

Many theoretical studies have been performed to obtain a molecular picture explaining the observed ultrafast decay mechanisms. Reaction paths and conical intersections between energy surfaces of different electronic states have been computed for the natural nucleobases,<sup>12–24</sup> nucleobase models,<sup>25–27</sup> isomers,<sup>17,28</sup> and substituted species.<sup>7,20,23,26</sup> Special emphasis must be attributed to these conical intersections and their energetic accessibility since they are responsible for the ultrafast nonradiative decay occurring in the nucleobases.

Concentrating the discussion on uracil, the energetically lowest conical intersection found so far on the  $S_2/S_1$  crossing seam has a boat conformation.<sup>19,29,30</sup> Two main conical intersections for  $S_1 \rightarrow S_0$  deactivation were identified for uracil. The first one is an ethylenic conical intersection formed by twisting of the  $C_5-C_6$  bond, which promotes a crossing between the  $\pi\pi^*$  and closed shell (cs) states. The second conical intersection, found by Lan et al. using the semiempirical OM2 approach, is formed by out-of-plane distortion of the oxygen atom connected to the  $C_4$  atom, causing a crossing between the  $n\pi^*$  and the closed shell states.<sup>31</sup> On the basis of these investigations, different mechanisms have been proposed to explain the photophysical deactivation to the

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## The decay mechanism of photoexcited guanine – A nonadiabatic dynamics study

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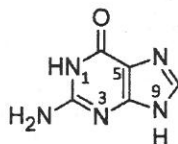
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*Ab initio* surface hopping dynamics calculations were performed for the biologically relevant tautomer of guanine in gas phase excited into the first  $\pi\pi^*$  state. The results show that the complete population of UV-excited molecules returns to the ground state following an exponential decay within  $\sim 220$  fs. This value is in good agreement with the experimentally obtained decay times of 148 and 360 fs. No fraction of the population remains trapped in the excited states. The internal conversion occurs in the  $\pi\pi^*$  state at two related types of conical intersections strongly puckered at the C2 atom. Only a small population of about 5% following an alternative pathway via a  $n\pi^*$  state was found in the dynamics. © 2011 American Institute of Physics. [doi:10.1063/1.3521498]

### I. INTRODUCTION

Like the four other natural nucleobases, UV-excited guanine shows ultrafast radiationless deactivation in gas phase with a decay time of a few picoseconds in the maximum,<sup>1–3</sup> which implies that it returns to the ground state by means of an internal conversion mechanism. These ultrafast deactivation mechanisms of guanine and of the other nucleobases and their related photostability are of great interest since they are considered to contribute to a natural chemical defense of the genetic code against damaging photochemically induced processes in reactive excited states.<sup>4</sup>

The guanine deactivation processes have been measured at 267 nm pump wavelength,<sup>1–3</sup> slightly below the center of the first absorption band (284 nm).<sup>5</sup> The exact time constants depend on the details of the experimental setup and on the way of the transient spectra are analyzed. When only the non-fragmented guanine channel is included in the analysis of the transient spectrum, guanine deactivation has been explained by two time constants of 0.148 and 0.36 ps.<sup>2</sup> When, however, fragment mass channels are included as well, the longer time constant has been found to increase more than six times to 2.3 ps.<sup>3</sup>



SCHEME 1.

Three main reaction pathways have been identified computationally for internal conversion of guanine (9H-1H-keto-

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guanine, Scheme 1) from the first  $\pi\pi^*$  state to the ground state. The first one is a barrierless pathway with the electronic state characterized as  $\pi\pi^*$  excitation and the nuclear motion described by puckering of the C2 atom and displacement of the amino group out of the ring plane.<sup>6–10</sup> Two slightly different kinds of conical intersections (denoted *ethylenic I* and *ethylenic II*) can be accessed through this pathway. The second pathway leads to a conical intersection characterized by puckering of the C6 atom and displacement of the oxygen atom out of the ring plane.<sup>6–8</sup> The third alternative is characterized by stretching of one of the NH bonds, promoting an intersection of the  $\pi\sigma^*$  state with the ground state. Reference 9 shows that this pathway may contribute to the deactivation for excitation energies around 0.4 eV above the absorption maximum of the first bright  $\pi\pi^*$  state. In the present, work we concentrated on a narrower range of excitation energies below values relevant to N-H dissociation and, therefore, did not consider this reaction channel further.

The guanine decay dynamics (see Ref. 2 and discussion above) has been interpreted by Serrano-Andrés *et al.*<sup>7</sup> by a three-state model based on computed minimum energy paths. The fast decay time observed experimentally (0.148 fs) has been assigned to a direct decay path to the first type of conical intersection (*ethylenic II*), whereas for the explanation of the longer decay time (360 fs) another alternative passing through the  $S_1$   $n\pi^*$  minimum has been suggested. Such specific assignments based on static calculations alone are difficult. Dynamics simulations are certainly the appropriate tool to resolve such mechanistic alternatives directly.

So far, non-adiabatic dynamics simulations for excited guanine have been performed with potential energy surfaces computed only with restricted open-shell Kohn–Sham (ROKS)<sup>11</sup> and semi-empirical OM2/MRCI<sup>8</sup> methods, predicting strongly differing results. The ROKS simulations for 9-methyl-guanine and 7-methyl-guanine do not show the involvement of conical intersections. Instead, the internal

**PHOTODYNAMICS OF THE ADENINE MODEL 4-AMINOPYRIMIDINE EMBEDDED WITHIN DOUBLE STRAND OF DNA**Tomáš ZELENÝ<sup>a</sup>, Pavel HOBZA<sup>b1</sup>, Dana NACHTIGALLOVÁ<sup>b2,\*</sup>,  
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*Dedicated to Dr. Zdeněk Havlas on the occasion of his 60th birthday.*

On-the-fly surface hopping nonadiabatic photodynamical simulations using hybrid quantum mechanical/molecular mechanical approach of 4-aminopyrimidine were performed to model the relaxation mechanism of adenine within DNA double strand. The surrounding bases do not affect the overall ring-puckering relaxation mechanisms significantly, however, interesting hydrogen-bond dynamics is observed. First, formation of intra-strand hydrogen bonds is found. It is shown that this effect speeds up the decay process. In addition, the Watson-Crick structure is altered by breaking one of the inter-strand hydrogen bonds also leading to a decrease of the life time.

**Keywords:** *Ab initio* calculations; Excited states; Nucleic acids; Photodynamics; QM/MM method.

The nature of the excited states of DNA/RNA nucleic acids induced by UV radiation has been studied extensively over the last decades. To understand the photochemistry of these compounds, the excited state behavior of individual bases<sup>1-19</sup>, base pairs<sup>20-24</sup> and nucleic acid strands<sup>25-31</sup> has been studied both experimentally and theoretically. The ultrafast relaxation of excited states of isolated species, on the time scale of a few picoseconds has been the subject of several computational studies performed by means of nonadiabatic dynamics simulations<sup>11-15,32,33</sup>. It is now generally accepted that nucleobases relax into the ground state through nonadiabatic transi-

## Does Stacking Restrain the Photodynamics of Individual Nucleobases?

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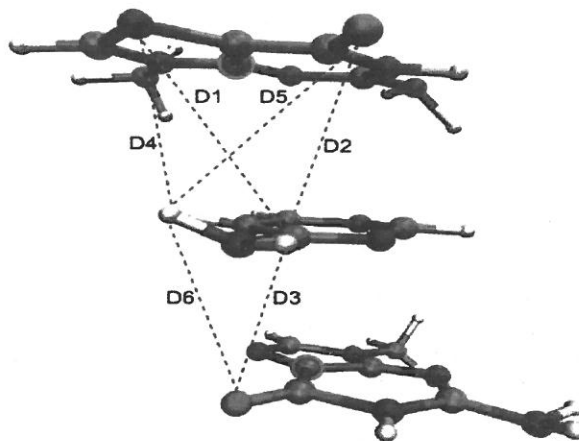
Nonadiabatic photodynamical simulations of 4-aminopyrimidine (4-APy) used as a model for adenine were performed by embedding it between two stacking methyl-guanine (mGua) molecules to determine the effect of spatial restrictions on the ultrafast photo-deactivation mechanism of this nucleobase. During the dynamics the formation of a significant fraction of intrastrand hydrogen bonding from 4-APy to mGua above and below is observed. These findings show that this type of hydrogen bond may play an important role for the photodynamics within one DNA strand and that it should be of interest even in irregular segments of double stranded nucleic acids structures. The relaxation mechanism of internal conversion to the ground state is dominated by ring puckering, and an overall elongation of the lifetime by ~20% as compared to the isolated 4-APy is computed.

Ultrafast relaxation of excited states of naturally occurring nucleobases on the time scale of a few picoseconds is a well-known phenomenon.<sup>1–4</sup> Computational studies performed by means of static<sup>5–8</sup> and nonadiabatic dynamics simulations<sup>9–12</sup> helped to reveal the excited-state time evolution of nucleobases as isolated species. It is generally accepted that isolated nucleobases relax into the ground state through nonadiabatic transitions via conical intersections characterized by strongly ring-puckered structures. Experimental observations<sup>13–15</sup> supported by theoretical calculations<sup>16–18</sup> indicate that in DNA excimers and/or excitons can be formed affecting the resulting excited-state lifetime strongly. Additionally, localization of the excitation on a single base, which is expected to resemble the behavior of the isolated base, has been reported as well.<sup>14</sup>

The theoretical treatment of the photodynamics of DNA, especially explicit nonadiabatic simulations, is quite involved, and corresponding investigations are only at the beginning.<sup>19–21</sup> In this contribution we focus in a first step on steric effects which adjoining nucleobases will have on an embedded base and its possibilities to reach conical intersections. These calculations are performed in the framework of on-the-fly surface hopping dynamics simulations which provide a detailed picture of the occurring photophysical processes. A combined quantum mechanical/molecular mechanical (QM/MM)<sup>22–24</sup> approach is used in which 4-APy is embedded between two methyl-guanine molecules (see Figure 1). The methyl groups were used to terminate the N-glycosidic bond in Gua. In these calculations the electronic excitation is confined to 4-APy, which is treated quantum mechanically, whereas the mGua mol-

ecules are treated at the MM level. Although this approach does not account for the electronic reactivity of 4-APy with adjacent nucleobases, it provides information on their possible sterical constraints during the excited state decay. Based on previous experience with dynamics simulations on isolated 4-APy<sup>25–27</sup> this molecule has been chosen because it is flexible enough to allow investigating restrictions to puckering at C<sub>2</sub> and C<sub>4</sub> atoms, relevant for purine bases, and at the C<sub>6</sub> atom, relevant for pyrimidine bases (for numbering scheme, see Figure 2).

The stacking model of mGua–4-APy–mGua was obtained by replacing adenine with 4-APy in the Gua–Ade–Gua sequence cut out from a B-DNA dodecamer crystal structure (PDBID: 196D).<sup>28</sup> These results are compared to those obtained from photodynamics simulation of isolated 4-APy<sup>27</sup> performed with identical procedures for creating initial conditions as used in the present study.



**Figure 1.** Starting structure of mGua–4-APy–mGua for dynamics calculations. Distances D1–D6 were used to analyze possible hydrogen bonds forming through the dynamics simulation. The atoms constrained by harmonic potential function are highlighted by red circles.

The 4-APy (QM part) was calculated at the state-averaged complete active space self-consistent field (SA-CASSCF) level of theory. A state averaging procedure with three states (SA-3) was used throughout the calculations. The active space was composed of two lone pair orbitals localized on nitrogen atoms of the pyrimidine ring and three  $\pi$  orbitals and three  $\pi^*$  orbitals, i.e., 10 electrons in 8 orbitals (CAS(10,8)). The molecular orbitals were expanded in the 6-31G\* basis set.<sup>29,30</sup> Analytic energy gradients and nonadiabatic coupling vectors required for the dynamics were computed as described in refs 31–34 using the COLUMBUS program system.<sup>35–37</sup>

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## Strikingly Different Effects of Hydrogen Bonding on the Photodynamics of Individual Nucleobases in DNA: Comparison of Guanine and Cytosine

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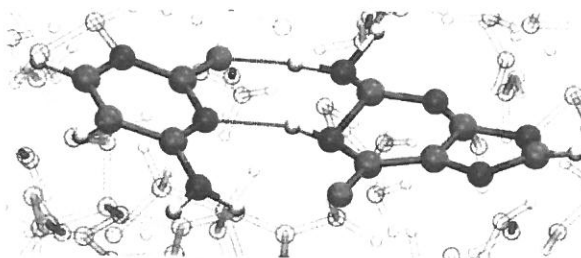
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### Supporting Information

**ABSTRACT:** Ab initio surface hopping dynamics calculations were performed to study the photophysical behavior of cytosine and guanine embedded in DNA using a quantum mechanical/molecular mechanics (QM/MM) approach. It was found that the decay rates of photo excited cytosine and guanine were affected in a completely different way by the hydrogen bonding to the DNA environment. In case of cytosine, the geometrical restrictions exerted by the hydrogen bonds did not influence the relaxation time of cytosine significantly due to the generally small cytosine ring puckering required to access the crossing region between excited and ground state. On the contrary, the presence of hydrogen bonds significantly altered the photodynamics of guanine. The analysis of the dynamics indicates that the major contribution to the lifetime changes comes from the interstrand hydrogen bonds. These bonds considerably restricted the out-of-plane motions of the NH<sub>2</sub> group of guanine which are necessary for the ultrafast decay to the ground state. As a result, only a negligible amount of trajectories decayed into the ground state for guanine embedded in DNA within the simulation time of 0.5 ps, while for comparison, the isolated guanine relaxed to the ground state with a lifetime of about 0.22 ps. These examples show that, in addition to phenomena related to electronic interactions between nucleobases, there also exist relatively simple mechanisms in DNA by which the lifetime of a nucleobase is significantly enhanced as compared to the gas phase.



### INTRODUCTION

In the last decades, impressive progress has been achieved in the understanding of the photodynamics of nucleic acids. The experimentally observed ultrafast excited state relaxation of all five naturally occurring nucleobases<sup>1–6</sup> has been explained by internal conversion via conical intersections.<sup>7–18</sup> The excited state behavior of nucleobases embedded within the DNA or RNA structure is, however, much less understood. The complex behavior of their excited state decay made it difficult to resolve the relaxation mechanisms in molecular detail. It has been shown in several experimental studies that the decay properties of nucleic acids in DNA depend on the base sequence and the conformation.<sup>4,19–28</sup> In contrast to isolated nucleic acid bases in gas phase and water, where conversion to the ground state occurs on picosecond or subpicosecond time scale, the relaxation of nucleobases in DNA ranges in the picosecond

and up to hundreds of picoseconds time scale. On the basis of the experimental observations, it has been suggested that in addition to excited states localized on single nucleobases also delocalized exciton and excimer/excimer complexes are responsible for nonradiative relaxation of nucleic acids.<sup>19,21,23,29–38</sup> Static quantum chemical calculations which can provide information about the character of excimer/excimer complexes and the influence of both, inter- and intrastrand interactions within the DNA helix<sup>39–48</sup> support these suggestions. The relaxation mediated by a proton-transfer within base-pairs was suggested also.<sup>49</sup> The role of the second strand, and consequently this type of relaxation mechanism, was, however, lately questioned by Crespo-Hernandez et al.<sup>19</sup>

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## The photodynamics of 2,4-diaminopyrimidine in comparison with 4-aminopyrimidine: The effect of amino-substitution

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### ABSTRACT

Nonadiabatic photodynamical surface hopping simulations of 4-aminopyrimidine and 2,4-diaminopyrimidine were performed at multi-configurational level. Additional NH<sub>2</sub> substitution blocks part of the ring puckering modes necessary to reach ultrafast deactivation channels but does not affect the excited-state lifetime of 2,4-diaminopyrimidine significantly since other pathways leading to ultrafast relaxation can be utilized.

The effect of the excess energy on the relaxation mechanism of 2,4-diaminopyrimidine was studied. It is shown that the amount of initial energy influences the time to reach the crossing seam and thus the lifetime. This study is important when interpreting the experimental results performed at different pump energies.

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### 1. Introduction

The nature of the excited state of DNA and RNA bases have been recently studied both experimentally and theoretically by several groups. In the gas phase the five naturally occurring nucleic acid bases have an ultrafast relaxation of the excited state to the ground state [1–6] which is explained by a rapid internal conversion via conical intersections [7–19].

Structural differences are known to modify the excited-state lifetime of these species and their related compounds. Long lived excited states of other tautomeric forms than naturally occurring nucleic acid bases have been reported [7,20]. The substitution of certain parts of the ring might affect the excited states lifetimes by blocking these parts and making relevant conical intersections inaccessible. A well known example is 2-aminopurine whose amino group at the C2 position leads to a much longer excited-state lifetime than of adenine (6-aminopurine) [21,22].

Study of the photophysics and photochemistry of model compounds with substitutions at the different position of the rings is especially important for the understanding of the nature of their ultrafast deactivation through specific conical intersection. One of the main theoretical tools in these investigations is the nonadiabatic dynamics simulation which can provide information about

the main relaxation pathways and the time scales of the relaxation process.

Recently, we have performed a combined experimental and theoretical study on the modified DNA base 2,4-diaminopyrimidine (2,4-DAPy; see Scheme 1) [23].

The static quantum chemical calculations clearly show that conical intersections are accessible when this molecule is excited near the absorption band maximum and, thus, an ultrafast (few ps) radiationless decay was predicted. In the experiments performed in the just-mentioned work pump energies near to the band origin were used revealing a much longer (estimated in the time range of 10 ps–1 ns) excited-state lifetime than what would be expected for a radiationless decay. This situation implies that sufficient energy must be available to overcome barriers before entering into the internal conversion regime and the photodynamics should be sensitive in this case to the pump energy used.

It was the goal of this work to study the photophysics of 2,4-DAPy by means of dynamics simulations in order to obtain explicit information on reaction mechanisms and lifetimes. The first objective was the evaluation of the overall influence of substituents on the relaxation mechanism of the pyrimidine ring by comparing 2,4-DAPy with 4-APy. The additional NH<sub>2</sub> substitution at the C2 atom can block specific decay channels and, consequently, increase the excited-state lifetime of 2,4-DAPy compared to 4-APy. The second aspect of this investigation concerns the afore-mentioned effect of the excess energy on the character of the dynamics by monitoring the dynamics simulations in dependence of different excitation energies. With higher energies a larger variety of conical intersections located on different crossing seams

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# Effect of substituents on the excited-state dynamics of the modified DNA bases 2,4-diaminopyrimidine and 2,6-diaminopurine†

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To explore the excited state dynamics of pyrimidine derivatives, we performed a combined experimental and theoretical study. We present resonant two-photon ionization (R2PI) and IR-UV double resonance spectra of 2,4-diaminopyrimidine and 2,6-diaminopurine seeded in a supersonic jet by laser desorption. For 2,4-diaminopyrimidine ( $S_0 \rightarrow S_1$  34 459  $\text{cm}^{-1}$ ), we observed only the diamino tautomer with an excited state lifetime bracketed between experimental limits of 10 ps and 1 ns. For 2,6-diaminopurine, we observed two tautomers, the 9H- ( $S_0 \rightarrow S_1$  34 881  $\text{cm}^{-1}$ ) and 7H- ( $S_0 \rightarrow S_1$  32 215  $\text{cm}^{-1}$ ) diamino forms, with excited state lifetimes of  $6.3 \pm 0.4$  ns and  $8.7 \pm 0.8$  ns, respectively. We investigated the nature of the excited state of 2,4-diaminopyrimidine by means of multi-reference *ab initio* methods. The calculations of stationary points in the ground and excited states, minima on the  $S_0/S_1$  crossing seam and connecting reaction paths show that several paths with negligible barriers exist, allowing ultrafast radiationless deactivation if excited at energies slightly higher than the band origin. The sub-nanosecond lifetime found experimentally is in good agreement with this finding.

## Introduction

Gas phase laser spectroscopy provides the means to study the intrinsic properties of biologically relevant molecules in isolation. Such studies on RNA and DNA bases have revealed unique photophysical properties that are sensitive to subtle structural differences. In many cases, the biologically most relevant tautomeric form has a sub picosecond excited state lifetime, while other *tautomeric forms of the same compound* are much longer lived.<sup>1,2</sup> We have even found that the Watson–Crick structure, adopted by the guanine–cytosine

(GC) base pair in DNA, appears to have a much shorter excited state lifetime than other *structures of the same base pair*.<sup>3</sup> Theoretical models explain these short lifetimes by a rapid internal conversion in which the excited state ( $S_1$ ) is coupled to the ground state ( $S_0$ ) *via* pathways with no or a very small barrier leading to conical intersections.<sup>4</sup> For the longer lived structures, small differences in relative energies cause the existence of barriers that lead to discrete spectra and lifetimes that can be two orders of magnitude longer. This rapid internal conversion pathway provides selected isomers with significantly enhanced photochemical stability, absent in the other longer lived structures. It is conceivable that these differences between excited state lifetimes of different bases and base-pair structures could have played a significant role in prebiotic chemistry.<sup>3</sup>

The five naturally occurring nucleic acid bases<sup>5–9</sup> exhibit an ultrafast excited state relaxation in the gas phase. For DNA bases it has been shown<sup>1,10–24</sup> that the energetically lowest conical intersections responsible for fast deactivation to the ground state are characterized by ring puckering modes. The most favorable ring puckering conical intersections result from the change of the HC6C5R (R = H, CH<sub>3</sub>) dihedral angle in uracil and thymine and the HNIC2H dihedral angle in adenine. This model also explains why by contrast 2-aminopurine, in which the C2 position is substituted by the amino group, has a long excited state lifetime and strong fluorescence.<sup>25–27</sup> Alternatively, the puckering at the C6 atom was also suggested to be responsible for adenine relaxation.<sup>28</sup>

Recently, surface-hopping dynamics studies of 4-aminopyrimidine (4-APy)<sup>29–31</sup> and 9H-adenine<sup>24</sup> showed that while the former relaxes into the ground state *via* different conical

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† Electronic supplementary information (ESI) available: The calculated vertical ionization potentials for all the 2,4-diaminopyrimidine tautomers (Fig. S1 and Table S1). The vertical excitation energies calculated at the RICC2 method with various basis sets (Table S2). The interpolation curves between the S1min\_C2 and S1min\_C6 minima towards various MXS structures calculated at the CASSCF and MR-CISD(17)+Q methods (Fig. S2). This material is available free of charge on the Internet. See DOI: 10.1039/b917852j

# Electronic splitting in the excited states of DNA base homodimers and -trimers: an evaluation of short-range and Coulombic interactions†

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The nature of the electronic interactions of the stacked nucleic-acid bases (adenine, thymine, cytosine, and uracil) in homodimer and -trimer complexes in their electronically excited states was investigated and analysed in terms of orbital-overlap and Coulombic interactions. The mutual orientations of the adjacent bases were selected so as to correspond to the A- and B-DNA conformations. The extent of the electronic interaction is larger for the former conformation. It is shown that the orbital-overlap interactions at the distance of two bases relevant to the DNA structure do not contribute significantly to the overall electronic coupling. Only the states which are caused by the ( $\pi \rightarrow \pi^*$ ) transitions manifest an electronic coupling.

## 1. Introduction

The phenomenon of the delocalisation of excited states, the electronic coupling and the transfer of excitation energy between the adjacent nucleic-acid bases in the stacked conformation has recently become the subject of many discussions. These processes are closely related to energy migration along the nucleic-acid strands, which may be a possible mechanism through which nucleic acids protect themselves against UV damage. The behaviour of nucleic acids in the excited states is very complex. The dynamics of nucleic acids in the excited states are affected by numerous factors, including *e.g.* the plasticity of the sugar–phosphate backbone, which influences the mutual orientation of bases, the lifetime of the excited states and the presence of water molecules, which are bonded directly to the nucleic-acid bases. It is thus difficult to identify the extent of the electronic coupling between the stacked nucleic-acid bases, which is responsible for both the formation of excimers and the energy migration.

The development of time-resolved fluorescence spectroscopy with femtosecond resolution makes it possible to investigate the excited states of nucleic-acid bases and the nature of their interactions. The features observed in the spectra can be explained by the presence of excited states localised on monomers, excimer formation and the existence of excitons. The effects of base sequence, the size of the DNA duplex and the conformation on the spectra were examined.<sup>1–8</sup> Spectral manifestation of the exciton formation was observed for adenine and cytosine dimers but not for thymine and guanine dimers.<sup>8</sup> Similar differences with respect to excimer

formation were observed also for adenine and cytosine, as well as thymine and guanine oligonucleotides.<sup>9</sup>

Experiments performed with the alternating (dA·dT), (dG·dC), and non-alternating dA and dT polymers or oligomers have shown that the efficiency of excimer formation depends on the base sequence.<sup>3,10–12</sup> In particular, the transient absorption signals associated with the excimer formation decay faster in the case of oligomers with (dA·dT) alternating sequences. The slow component of the transient signals of non-alternating sequences in (dA)<sub>18</sub>·(dT)<sub>18</sub> was very similar to that observed for (dA)<sub>18</sub>. It was suggested that the excited states are initially localised, with the subsequent ultrafast formation of excimers.<sup>3</sup>

Based on the steady-state absorption and fluorescence spectra of poly(dA)·poly(dT), it was suggested that the excited states are delocalised over several bases.<sup>4,13</sup> Onidas *et al.* have studied the fluorescence spectra of alternating A–T polymers and oligomers and compared them to those for the homopolymeric sequences (dA)<sub>n</sub>·(dT)<sub>n</sub>.<sup>14</sup> The steady-state fluorescence spectra for these two types of oligomers and polymers were different. The differences in the spectra were explained by a population of exciton states after the absorption, followed by an ultrafast intra-band scattering and emission of monomers and excitons. In the case of the alternating (dAdT)<sub>n</sub>·(dAdT)<sub>n</sub>, these emitting states were quenched by excimers characterised by long-lived excited states. The nature of the energy transfer was investigated theoretically for the model B-DNA poly(dA)·poly(dT) duplex.<sup>15</sup> In accord with experimental findings, it was suggested that for exciton dynamics the intra-strand stacking interactions are much more important than those between two chains. The calculations also predict that, unlike the exciton located on adenosine, the thymidine exciton decomposes into electron–hole pairs.

Miannay *et al.* have recently reported the results of a femtosecond study of guanine–cytosine DNA double helices.<sup>5</sup> In poly(dGdC)·poly(dGdC), excited states relax faster than the respective monomers, unlike A–T duplexes, characterised by a lengthening of the excited-state lifetime. Nevertheless, signs of the electron transfer along the helix have been detected also for this system.

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† Electronic supplementary information (ESI) available: Vertical excitation energies and details of various approximations for the transition Coulomb interaction. See DOI: 10.1039/b806323k



# Electronic coupling in the excited electronic state of stacked DNA base homodimers†

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The nature of the electronic coupling of stacked nucleic acid bases adenine (A), thymine (T), and cytosine (C), in A–A, T–T, and C–C complexes in their excited states was investigated; a different character of the electronic coupling for the T–T complex was shown.

The nature of the excited states of DNA bases, in particular, the extent of excited state delocalization along the double helix has been the subject of much discussion. The lack of any significant spectral shift in UV spectra<sup>1</sup> led to the opinion that the excited states are localized on a single base contrary to theoretically suggested delocalization of the excited states.<sup>2</sup> Only recently has the use of time-resolved fluorescence spectroscopy with femtosecond resolutions of DNA oligomers<sup>3–11</sup> allowed the reinvestigation of the problem of excited states delocalization. Based on the results of these studies, in particular, on single and double helices of A and T, it is now assumed that the excited states are delocalized over several bases. Slow components which appear in time-resolved transient absorption spectra of (dA)<sub>18</sub>, (dA)<sub>18</sub>·(dT)<sub>18</sub>, and (dAdT)<sub>9</sub>·(dAdT)<sub>9</sub> were associated with the adenine excimers.<sup>3</sup>

The question of excited states delocalization becomes important when discussing the photostability of DNA with respect to UV damage. One of the suggested operative mechanisms is linked to the very short lifetime of excited states of isolated nucleic acid bases<sup>12</sup> which is attributed to non-radiative relaxation to the ground state. Several mechanisms were suggested for this non-radiative decay.<sup>13–17</sup>

An alternative explanation of how DNA protects itself from UV damage is the fast transfer of excitation energy which prevents the localization of electronic energy and subsequent photochemical reaction. It has already been suggested that the differences in the photochemistry of poly(A) and poly(dA) homopolymers are the result of different stacking of the two

adjacent bases.<sup>18</sup> While poly(A) which possesses A-type DNA conformation shows almost no reactivity in the excited state, poly(dA) with B-type DNA conformation is able to form dimeric photoproducts. The photodimerization process of two adjacent pyrimidine bases is a problem which attracts even more attention. There are two experimental observations which are subject to theoretical investigation: (i) upon UVC irradiation *cis,syn*-cyclobutane dimer is most frequently formed between two adjacent T bases,<sup>19</sup> (ii) oligonucleotides in an A-type conformation are more UV resistant than in B-type conformation. Also in this case the question of the role of delocalization was raised.<sup>5</sup> The mechanism of the photodimerization process of adjacent thymines was investigated theoretically using a time-dependent DFT method.<sup>20–23</sup>

Due to its structure, the electronic coupling and hence the nature of excited state delocalization results from both the orbital overlap in the short-range limit and interaction of transition dipole moments (dipole–dipole interaction) in the intermediate-range limit.<sup>24</sup> It is a generally accepted view that the extent of the delocalization depends on the base sequence as well as on the structure of the DNA helix which effects their mutual orientation.

Clearly, electron delocalization is a very complex problem in which several aspects, including the base sequence and respective orientation of the bases controlled by the sugar–phosphate structure play an important role in both the short-range and intermediate-range limits. The nature of excited states within the DNA helix has been studied theoretically within the framework of the exciton theory.<sup>8–10</sup> In these studies, the electronic coupling due to the dipole–dipole interaction only was investigated. In agreement with experiment, it was suggested that the excited states are delocalized over at least two bases. Very recently, Olaso-González *et al.* using *ab initio* complete active space second order perturbation theory (CASPT2) studied the lowest excited state of two cytosine bases in the stacked conformation.<sup>25</sup> The excimer formation was predicted which is suggested to be responsible for redshifted fluorescence observed in cytosine oligonucleotides.<sup>26</sup>

When two chromophores interact in the excited state, the resulting transition energies are shifted with respect to the energy of the monomer and the transition energies split. For symmetrical arrangements this splitting is a measure of the electronic coupling between the two chromophores (bases) in the excited state.<sup>24</sup>

For non-symmetrical arrangements of the monomers, the resulting energy splitting is due to both, a net electronic

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† Electronic supplementary information (ESI) available: Geometries of stacked DNA bases (Fig. S1–S3); excitation energies (in eV) and their splitting ( $\Delta E$ ) of model ethylene–dimer and ethylene–acetylene systems in stacked conformation (Tables S1 and S2). See DOI: 10.1039/b618382d

## PAPER

## The effect of dimerization on the excited state behavior of methylated xanthine derivatives: a computational study†

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The behavior of monomers and dimers of methylated xanthine derivatives in their excited states is investigated by means of the ADC(2), CASSCF, and CASPT2 methods. The results of the calculations of stationary points in the ground and excited states, minima on the  $S_0/S_1$  crossing seams and the relaxation pathways are used to provide the interpretation of experimental observations of the monomer xanthine derivatives. The effect of dimerization on the excited state properties is studied for various relative orientations of the monomers in the dimer complexes in comparison with the relevant monomer species. A significant stabilization in the excited state minima of dimers is observed. These can act as trapping sites. Various types of conical intersections, with both localized and delocalized characters of wavefunctions, have been found, mainly energetically above the lowest bright excited state in the FC region. In addition, structures with the bonds formed between the two monomers were also found on the crossing seams. The possibility of ultrafast relaxation *via* these conical intersections is discussed.

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## Introduction

In the last decade, impressive progress in the field of photochemistry of nucleic acids and related molecular systems could be observed. As UV irradiation can cause serious damage to DNA, leading to modifications in the genetic code, this topic is highly important. The picosecond and subpicosecond excited state photodynamics of isolated naturally occurring nucleobases was investigated experimentally<sup>1–6</sup> and their relaxation mechanisms were explained by a nonadiabatic internal conversion *via* conical intersections.<sup>7–19</sup> These processes are fairly well understood now, even though many interesting detailed questions still await their solution. The excited state behavior of nucleobases becomes much more complex when embedded in the structure of nucleic acids. The relaxation occurs on a time scale of picoseconds up to hundreds of picoseconds. Relaxation mechanisms mediated by both, the

formation of intra-strand delocalized excitons and excimer/excimer complexes, including charge transfer excited states,<sup>20–28</sup> and interstrand proton transfer, were discussed.<sup>29–32</sup>

To reduce this complexity, the photodynamics of dinucleotide analogues has been investigated (*e.g.* ref. 33) with the goal of transferring the experience obtained with such simpler problems to the more complex situation found in DNA. In this regard, nucleic acid analogues play an important role because of the larger choice of compounds available and the interest in replacing natural counterparts in DNA by these candidates. Xanthine (3,7-dihydro-purine-2,6-dione) is one important representative of such nucleic acid analogues.

Recently, excited state lifetimes of xanthine (X) and methylated derivatives (*theophylline* (1,3-diMet-X), *theobromine* (3,7-diMet-X) and *caffeine* (1,3,7-triMet-X)) were determined using femtosecond transient absorption spectroscopy in aqueous and acetonitrile solutions.<sup>34</sup> The observed spectra revealed that in water all of these species relax to the electronic ground state on a subpicosecond time scale. Resonant two-photon ionization (R2PI) and IR-UV double resonance spectroscopy performed for the same methylated species in the gas phase<sup>33</sup> and xanthine itself<sup>35</sup> were all sharp and vibrationally resolved, suggesting a long excited state lifetime. Computational studies based on the complete active space with perturbation theory to the second order (CASPT2) and approximate coupled cluster singles-and-doubles method (CC2) calculations assigned the gas phase monomer spectra to the 7*H-diketo* tautomer of

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# The Charge-Transfer States in a Stacked Nucleobase Dimer Complex: A Benchmark Study

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**Abstract:** Electronic singlet excitations of stacked adenine–thymine (AT) and guanine–cytosine (GC) complexes have been investigated with respect to local excitation and charge-transfer (CT) characters. Potential energy curves for rigid displacement of the nucleobases have been computed to establish the distance dependence of the CT states. The second-order algebraic diagrammatic construction [ADC(2)] method served as reference approach for comparison to a selected set of density functionals used within the time-dependent density functional theory (TD-DFT). Particular attention was dedicated to the performance of the recently developed family of M06 functionals. The calculations for the stacked complexes show that at the ADC(2) level, the lowest CT state is  $S_6$  for the AT and as  $S_4$  for the GC pair. At the reference geometry, the actual charge transferred is found to be  $0.73 e$  for AT. In case of GC, this amount is much smaller ( $0.17 e$ ). With increasing separation of the two nucleobases, the CT state is strongly destabilized. The M06-2X version provides a relatively good reproduction of the ADC(2) results. It avoids the serious overstabilization and overcrowding of the spectrum found with the B3LYP functional. On the other hand, M06-HF destabilizes the CT state too strongly. TD-DFT/M06-2X calculations in solution (heptane, isoquinoline, and water) using the polarizable continuum model show a stabilization of the CT state and an increase in CT character with increasing polarity of the solvent.

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**Key words:** charge transfer; excited states; coupled cluster; DFT; stacked nucleobases

Additional Supporting Information may be found in the online version of this article.

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# Singlet $L_a$ and $L_b$ Bands for N-Acenes ( $N = 2-7$ ): A CASSCF/CASPT2 Study

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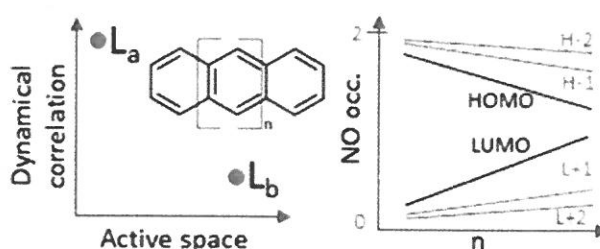
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## Supporting Information

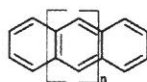
**ABSTRACT:** In this work CASPT2 calculations of polyacenes (from naphthalene to heptacene) were performed to find a methodology suitable for calculations of the absorption spectra, in particular of the  $L_a$  ( $B_{2u}$  state) and  $L_b$  ( $B_{3u}$  state) bands, of more extended systems. The effect of the extension of the active space and of freezing  $\sigma$  orbitals was investigated. The MCSCF excitation energy of the  $B_{2u}$  state is not sensitive to the size of the active space used. However, the CASPT2 results depend strongly on the amount of  $\sigma$  orbitals frozen reflecting the ionic character of the  $B_{2u}$  state. On the other hand, the excitation energies of the  $B_{3u}$  state are much more sensitive to the size of the active space used in the calculations reflecting its multiconfigurational character. We found a good agreement with experimental data for both bands by including 14 electrons in 14  $\pi$  orbitals in the active space followed by the CASPT2(14,14) perturbation scheme in which both  $\sigma$  and  $\pi$  orbitals are included.



## 1. INTRODUCTION

Due to the promising potential applications of materials which contain  $\pi$ -conjugated systems<sup>1-8</sup> to be used as chemical sensors, organic semiconductors, in spintronics and nonlinear optics, a large number of studies on graphene and its model systems appeared recently in the literature.<sup>9-15</sup> The electronic properties of polyacenes (PACs, Scheme 1) not only are very

Scheme 1. Structure of Studied Polyacenes ( $n = 0-5$ )



attractive for material science,<sup>16-24</sup> but also they can serve as a model system for studying the properties of ground and excited states of extended  $\pi$  systems<sup>23,25-42</sup> by means of theoretical methods.

PACs possess  $D_{2h}$  symmetry with a totally symmetric ( $A_{1g}$ ) electronic ground state. The two low-lying excited states in the absorption spectra are short-axis polarized ( $L_a$ ) and long-axis polarized ( $L_b$ ) states of  $B_{2u}$  and  $B_{3u}$  symmetries, respectively.<sup>33</sup> The two lowest excited states of gerade symmetry, not seen in

the absorption spectra, are of  $A_{1g}$  and  $B_{1g}$  symmetries. According to previous theoretical results,<sup>18,29,31-33,37-39,43-46</sup> two main conclusions can be drawn from these calculations: (1) the multireference character of the ground state electronic configuration increases with increasing number of fused benzene rings and (2) doubly excited configurations contribute to the wave functions of the low-lying excited states. For polyacenes larger than tetracene these configurations need to be included to correctly describe the ordering of the excited state of  $A_g$  and  $B_{3u}$  symmetries observed experimentally.<sup>44</sup>

A multiconfigurational character of the wave function was found<sup>37,43</sup> for the first excited state ( $B_{3u}$ ) of naphthalene, resulting from a linear combination of HOMO-1  $\rightarrow$  LUMO and HOMO  $\rightarrow$  LUMO+1 configurations, while the wave function of the second excited state ( $B_{2u}$ ) corresponds almost exclusively to a HOMO  $\rightarrow$  LUMO excitation. For acenes larger than naphthalene, the  $B_{2u}$  and  $B_{3u}$  states change their energy ordering. The  $A_{1g}$  state is best described by the double excitation HOMO<sup>2</sup>  $\rightarrow$  LUMO<sup>2</sup>. The  $B_{1g}$  state wave function has large contributions from the singly excited configurations

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# High-level theoretical benchmark investigations of the UV-vis absorption spectra of paradigmatic polycyclic aromatic hydrocarbons as models for graphene quantum dots

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## ABSTRACT

Five paradigmatic polycyclic aromatic hydrocarbons (PAHs) (pyrene, circum-1-pyrene, coronene, circum-1-coronene, and circum-2-coronene) are used for studying the performance of three single-reference methods [scaled opposite-spin-algebraic diagrammatic construction to second-order [SOS-ADC(2)], time-dependent (TD)-B3LYP, and TD-Coulomb-attenuating method (CAM)-B3LYP] and three multireference (MR) methods [density functional theory/multireference configuration interaction (DFT/MRCI), strongly contracted- $n$ -electron valence state perturbation theory to second order (NEVPT2), and spectroscopy oriented configuration interaction (SORCI)]. The performance of these methods was evaluated by comparison of the calculated vertical excitation energies with experiments, where available. DFT/MRCI performs best and thus was used as a benchmark for other approaches where experimental values were not available. Both TD-B3LYP and NEVPT2 agree well with the benchmark data. SORCI performs better for coronene than for pyrene. SOS-ADC(2) does reasonably well in terms of excitation energies for smaller systems, but the error increases somewhat as the size of the system gets bigger. The natural transition orbital analysis for SOS-ADC(2) results indicated that at least two configurations were essential to characterize most of the lower-case electronic states. TD-CAM-B3LYP gives the largest errors for excitation energies and also gives an incorrect order of the lowest two states in circum-1-pyrene. A strong density increase of dark states was observed in the UV spectra with increasing size except for the lowest few states which remained well separated. An extrapolation of the UV spectra to infinite PAH size for  $S_1$ ,  $S_2$ , and the first bright state based on the coronene series was made. The extrapolated excitation energies closest to experimental measurements on graphene quantum dots were obtained by TD-CAM-B3LYP.

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## I. INTRODUCTION

Polycyclic aromatic hydrocarbons (PAHs), conjugated aromatic molecules composed of fused benzenoid rings, have aroused widespread interest among scientists due to their important role

in many fields ranging from chemistry<sup>1</sup> and materials science<sup>2</sup> to astrochemistry.<sup>3</sup> In particular, PAHs are also regarded as a cut out and model of graphene, especially in the case of the larger PAHs denoted as nanographenes or graphene quantum dots (GQDs),<sup>4-6</sup> which are promising materials being used in

# Graphitic Nitrogen Triggers Red Fluorescence in Carbon Dots

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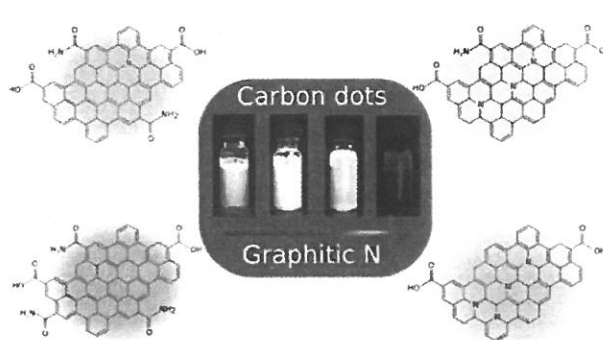
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## Supporting Information

**ABSTRACT:** Carbon dots (CDs) are a stable and highly biocompatible fluorescent material offering great application potential in cell labeling, optical imaging, LED diodes, and optoelectronic technologies. Because their emission wavelengths provide the best tissue penetration, red-emitting CDs are of particular interest for applications in biomedical technologies. Current synthetic strategies enabling red-shifted emission include increasing the CD particle size ( $sp^2$  domain) by a proper synthetic strategy and tuning the surface chemistry of CDs with suitable functional groups (e.g., carboxyl). Here we present an elegant route for preparing full-color CDs with well-controllable fluorescence at blue, green, yellow, or red wavelengths. The two-step procedure involves the synthesis of a full-color-emitting mixture of CDs from citric acid and urea in formamide followed by separation of the individual fluorescent fractions by column chromatography based on differences in CD charge. Red-emitting CDs, which had the most negative charge, were separated as the last fraction. The trend in the separation, surface charge, and red-shift of photoluminescence was caused by increasing amount of graphitic nitrogen in the CD structure, as was clearly proved by XPS, FT-IR, Raman spectroscopy, and DFT calculations. Importantly, graphitic nitrogen generates midgap states within the HOMO–LUMO gap of the undoped systems, resulting in significantly red-shifted light absorption that in turn gives rise to fluorescence at the low-energy end of the visible spectrum. The presented findings identify graphitic nitrogen as another crucial factor that can red-shift the CD photoluminescence.

**KEYWORDS:** nitrogen-doped, graphene dots, red fluorescence, fluorescence mechanism, band-gap tuning



Carbon dots (CDs) attract broad scientific interest owing to their attractive fluorescence properties. They are resistant to photobleaching akin to traditional semiconductor quantum dots<sup>1,2</sup> but still possess a great safety profile.<sup>3</sup> Therefore, they have found wide utilization in biomedical fields, e.g., cell labeling, optical imaging, drug delivery, and biosensing.<sup>4–9</sup> They have also been applied in various technologies, e.g., light-emitting diodes,<sup>10</sup> water splitting,<sup>11</sup> and photocatalysis.<sup>12</sup> In recent years, many synthetic procedures have been developed to prepare and tune the optical properties of CDs. Since the discovery of this material in 2004, mainly blue, green, and yellow fluorescent CDs have been reported.<sup>1,2</sup> However, very few reproducible synthetic strategies for red fluorescent CDs have been published.<sup>13,14</sup> Red-emitting

CDs are particularly suitable for biomedical applications because red light shows the best tissue penetration. In addition, red CDs can provide the missing building block for a full-color-emitting spectrum, e.g., for white-light-emitting diodes. Therefore, efficient and reproducible synthetic pathways for the production of red fluorescent CDs can advance applications of CDs in many scientific and industrial areas.

Recently, attention has been given to understanding the origin of CD photoluminescence because knowledge in this field may support the rational design of CD optical properties.

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## Progress and challenges in understanding of photoluminescence properties of carbon dots based on theoretical computations

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### ABSTRACT

Carbon dots (CDs), including graphene quantum dots, carbon nanodots, carbon quantum dots, and carbonized polymer dots, belong to extensively studied nanomaterials with a very broad application potential resulting from their bright photoluminescence (PL), high (photo)stability, low toxicity and great biocompatibility. However, the design of CDs with tailored properties is still hampered by a fairly limited understanding of the CD PL, which stems from their rather complex structure and variability of the PL centers. Theoretical calculations provide valuable insights into the nature of the excited states and the source of PL. In this review, we focus on state-of-the-art theoretical methods for the description of absorption and PL of CDs and their limitations, along with providing an overview of theoretical studies addressing structural models and the electronic structure of various types of CDs in the context of their overall optical properties. Besides the assessment of the current state of knowledge, we also highlight the opportunity for further advancements in the field.

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## 1. Introduction

### 1.1. Discovery and definition of carbon dots

Nanosized carbon-based materials, such as fullerenes, nanotubes, and graphene, have attracted attention of scientists owing to their unique physicochemical properties often very different from those of the original bulk material. In 2004, Xu *et al.* [1] started to write a new exciting chapter on carbon nanostructures by separating a mixture of weakly fluorescent carbon nanoparticles from single-walled carbon nanotubes obtained from arc-discharge soot. Since then, carbon dots (CDs), as these carbon-based fluorescent nanoparticles (with at least one dimension smaller than 10 nm) are nowadays referred to, have expanded into a large class of metal-free nanomaterials with significant and tuneable photoluminescence (PL). Other unique properties, in particular, non/low-toxicity, biocompatibility, dispersibility in a wide range of solvents, biodegradability, low-cost, facile and scalable synthesis, predispose CDs for plethora of environment-

friendly applications in biomedical imaging and sensing [2–4], cancer therapy [5], or theranostics [6], anti-counterfeiting [7], light-emitting devices [8–11], optoelectronic devices [12–15], photocatalysis [16] and functional materials [17,18].

### 1.2. Synthesis of CDs

In general, the CDs can be prepared by employing either top-down or bottom-up approaches [12]. The top-down methods usually require harsh conditions (strong oxidants, concentrated acids, and high temperatures) or severe physical techniques (laser ablation, arc discharge and nanolithography) to cut, decompose and/or exfoliate bulk carbonaceous materials (e.g., graphite, graphene oxide, activated carbon, soot, carbon nanotubes) into carbon nanoparticles [19–27]. On the other hand, the bottom-up strategies such as stepwise organic synthesis, hydrothermal/solvothermal synthesis, and chemical vapor deposition and microwave-assisted treatments can utilize a large variety of simple organic precursors [28,29]. The optimization of reagents and reaction conditions yielded CDs with significantly enhanced PL quantum yields (QYs) which were mainly attributed to heteroatom doping effects, surface functionalities, and, last but not least, the presence of molecular fluorophores

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# Tuning the UV spectrum of PAHs by means of different N-doping types taking pyrene as paradigmatic example: categorization *via* valence bond theory and high-level computational approaches†

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Tuning of the electronic spectra of carbon dots by means of inserting heteroatoms into the  $\pi$ -conjugated polycyclic aromatic hydrocarbon (PAH) system is a popular tool to achieve a broad range of absorption and emission frequencies. Especially nitrogen atoms have been used successfully for that purpose. Despite the significant progress achieved with these procedures, the prediction of specific shifts in the UV-vis spectra and the understanding of the electronic transitions is still a challenging task. In this work, high-level quantum chemical methods based on multireference (MR) and single-reference (SR) methods have been used to predict the effect of different nitrogen doping patterns inserted into the prototypical PAH pyrene on its absorption spectrum. Furthermore, a simple classification scheme based on valence bond (VB) theory and the Clar sextet rule in combination with the harmonic oscillator measure of aromaticity (HOMA) index was applied to arrange the different doping structures into groups and rationalize their electronic properties. The results show a wide variety of mostly redshifts in the spectra as compared to the pristine pyrene case. The most interesting doping structures with the largest red shifts leading to absorption energies below one eV could be readily explained by the occurrence of diradical VB structures in combination with Clar sextets. Moreover, analysis of the electronic transitions computed with MR methods showed that several of the low-lying excited states possess double-excitation character, which cannot be realized by the popular SR methods and, thus, are simply absent in the calculated spectra.

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## 1. Introduction

Due to their outstanding optical properties, carbon dots (CD) constitute a fascinating new carbon-based material with interesting applications in materials science, chemistry, biosensing and several other fields.<sup>1–5</sup> Their chemical modification, which

includes also chemical doping, provides a promising strategy to modulate their electronic properties which should convert them to materials with even better performance than their pristine counter-partners.<sup>6</sup> In particular, nitrogen doping changes the electronic properties of CDs, which can be efficiently used to design systems such as sensors,<sup>7</sup> electrocatalysts in redox reactions,<sup>8–12</sup> systems with capabilities of electrochemical energy storage,<sup>13</sup> photocatalysts,<sup>14–17</sup> and battery constructions.<sup>18</sup> Their optical properties with the photoluminescence quantum yield up to 75% make them efficient fluorescence probes<sup>19–23</sup> and light-emitting diodes.<sup>24</sup> In this context, the red-shift of the fluorescence is of great relevance *e.g.* for bioimaging applications because near infra-red light exhibits deeper tissue penetration.<sup>25</sup>

Despite the large potential of just-mentioned N-doped CDs, understanding of the optical spectra and finding materials which possess appropriately red-shifted UV and fluorescence properties represent still a challenging task for experimental as

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# Doping Capabilities of Fluorine on the UV Absorption and Emission Spectra of Pyrene-Based Graphene Quantum Dots

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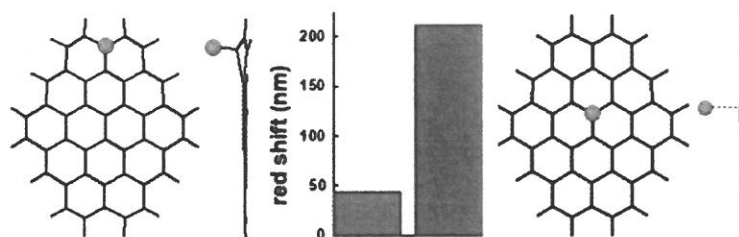
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**Fluorescence red-shift depends strongly on the F<sup>-</sup> position**

**ABSTRACT:** Functionalization of quantum carbon dots (QCDs) and graphene quantum dots (GQDs) is a popular way to tune their optical spectra increasing their potential applicability in material science and biorelated disciplines. Based on the experimental observation, functionalization by fluorine atoms induces substantial shifts in absorption and emission spectra and an intensity increase. Understanding of the effects due to fluorine functionalization at the atomic scale level is still challenging due to the complex structure of fluorinated QCDs. In this work, the effect of covalent edge-fluorination and fluorine anion doping on absorption and emission spectra of prototypical polycyclic aromatic hydrocarbons pyrene and circum-pyrene has been investigated. The ways to achieve efficient red-shifts in the UV spectra and obtaining reasonable intensities stood in the focus of the work. High-level quantum chemical methods based on density functional theory/multireference configuration interaction (DFT/MRCI) and single-reference second-order algebraic diagrammatic construction (ADC(2)) and density functional theory (DFT) using the CAM-B3LYP functional have been used for this purpose. The calculations show that doping with the fluoride anion can have significant effects on the electronic spectrum. However, the effect of the fluoride ion is strongly dependent on its position with respect to the QCD. The localization above the GQDs causes large red-shifts to both the absorption and emission of spectra of GQDs, while in-plane localization leads to only negligible shifts and a tendency to dissociation after electronic excitation. Thus, large red-shifts, observed in complexes with F<sup>-</sup>, are obtained due to the introduction of new excited states with large CT character not yet been considered previously in this context, although they have the potential to significantly influence the photophysics of quantum dots. Doping by edge fluorination redshifts the spectra only slightly. This study provides insights on fluorine-doped GQDs, which is conducive to promoting its rational design and controllable synthesis.

## 1. INTRODUCTION

Fluorination of carbon-based materials<sup>1</sup> is a widely investigated way to modify the properties of carbon materials, including graphite,<sup>2</sup> graphene,<sup>3–8</sup> and carbon nanotubes.<sup>2,9–11</sup> Extensive research on fluoro-doped carbon materials, performed over the last 10 years, proved their unique electronic,<sup>7,12</sup> magnetic,<sup>13,14</sup> and luminescent<sup>15,16</sup> properties. Their high reactivity, thermal stability, and hydrophobic surface properties make these compounds an extraordinary material for applications in electrocatalysis, electrochemical sensing, spintronics, anticorrosion, separation processes, optoelectronics and photonics, and biomedicine.<sup>8</sup> In particular, the red-shift of the fluorescence emission frequency is of great interest for bioimaging applications because near-infrared light exhibits deeper tissue

penetration.<sup>17</sup> The origin of the red-shifts, observed during fluorine doping of polyaromatic hydrocarbons dibenzoanthracene- and benzoperylene-type systems were explained by the electron-withdrawing character of fluorine.<sup>18</sup>

The stability of the C–F bond and the hydrophobic nature of the quantum carbon dots (QCD) make the synthesis of fluorinated quantum carbon dots (F-QCD) more challenging

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## Excited states and excitonic interactions in prototypic polycyclic aromatic hydrocarbon dimers as models for graphitic interactions in carbon dots†

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The study of electronically excited states of stacked polycyclic aromatic hydrocarbons (PAHs) is of great interest due to promising applications of these compounds as luminescent carbon nanomaterials such as graphene quantum dots (GQDs) and carbon dots (CDs). In this study, the excited states and excitonic interactions are described in detail based on four CD model dimer systems of pyrene, coronene, circum-1-pyrene and circum-1-coronene. Two multi-reference methods, DFT/MRCI and SC-NEVPT2, and two single-reference methods, ADC(2) and CAM-B3LYP, have been used for excited state calculations. The DFT/MRCI method has been used as a benchmark method to evaluate the performance of the other ones. All methods produce useful lists of excited states. However, an overestimation of excitation energies and an inverted ordering of states, especially concerning the bright HOMO–LUMO excitation, are observed. In the pyrene-based systems, the first bright state appears among the first four states, whereas the number of dark states is significantly larger for the coronene-based systems. Fluorescence emission properties are addressed by means of geometry optimization in the  $S_1$  state. The inter sheet distances for the  $S_1$  state decrease in comparison to the corresponding ground-state values. These reductions are largest for the pyrene dimer and decrease significantly for the larger dimers. Several minima have been determined on the  $S_1$  energy surface for most of the dimers. The largest variability in emission energies is found for the pyrene dimer, whereas in the other cases a more regular behavior of the emission spectra is observed.

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### 1 Introduction

Excited state aggregates of polycyclic aromatic species, which are the building blocks of graphene quantum dots (GQDs) and carbon dots (CDs), are currently extensively studied systems as promising luminescent carbon nanomaterials<sup>1</sup> due to their potential applications in photonics and optoelectronics, organic devices, singlet fission processes, and bio- and medical applications.<sup>2–7</sup> Understanding and controlling the origin of the photoluminescence

of GQDs and CDs, which can help in the rational design of carbon-based materials, was the subject of experimental and computational research (for recent studies see *e.g.* ref. 8–10). GQDs can be characterized as larger polycyclic aromatic hydrocarbons (PAHs), which obtain their discrete absorption and emission spectra due to their quantum confinement, as compared to graphene sheets.<sup>11</sup> CDs are more complex 3-dimensional systems which contain as important components graphitic regions responsible for their photoluminescence properties.<sup>12</sup> Due to the aggregation of PAHs in the graphitic regions of stacked sheets, the optical properties often resemble those of excimers (excited state dimers of two identical monomers) or exciplexes (excited state complexes of two non-identical monomers), complexes weakly associated in the ground state, but strongly interacting in the excited states.<sup>13,14</sup> Potential applications of these associates, including those in CDs, were already suggested in several studies<sup>15–21</sup> and their optical properties were investigated also computationally using model systems.<sup>22–24</sup>

The first experimentally observed pyrene excimer,<sup>25</sup> as one of the basic PAH excimer models, has received great attention

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# Emission Energies and Stokes Shifts for Single Polycyclic Aromatic Hydrocarbon Sheets in Comparison to the Effect of Excimer Formation

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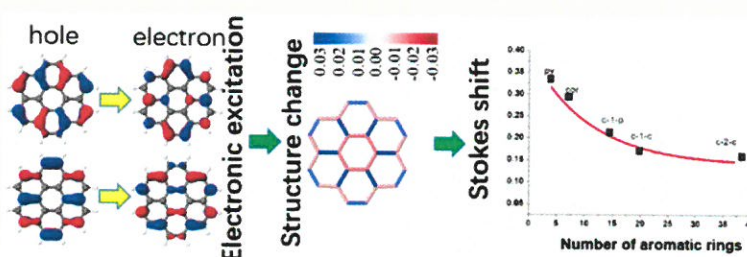
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## Supporting Information



**ABSTRACT:** Emission spectra of paradigmatic single-sheet polycyclic aromatic hydrocarbons (PAHs), pyrene, circum-1-pyrene, coronene, circum-1-coronene, and circum-2-coronene and Stokes shifts were computed and compared with previously calculated comparable data for relaxed excimer structures using the SOS-ADC(2), TD-B3LYP, and TD-CAM-B3LYP methods with multireference DFT/MRCI data as the benchmark. Vertical emission transitions and Stokes shifts were extrapolated to infinite PAH size. Comparison of Stokes shifts computed from theoretical monomer and dimer data confirms assumptions that relaxed excimers are responsible for the unusually large Stokes shifts in carbon dots observed in experimental investigations.

Polycyclic aromatic hydrocarbons (PAHs), conjugated aromatic molecules composed of fused benzenoid rings, have attracted attention as building units of carbon dot structures (CDs) which include graphene quantum dots (GQDs), composed of one or few graphene layers a few tens of nanometers in size, and carbon quantum dots (CQDs), possessing crystalline structure.<sup>1</sup> CDs are three-dimensional systems containing graphitic regions, which, in contrast to graphene, exhibit discrete absorption and emission spectra arising from their quantum confinement.<sup>1,2</sup> Because of their attractive photoluminescence properties and photostability, these materials find promising applications as photodetectors, solar cells, and light-emitting diodes.<sup>3,4</sup> The biocompatibility of CDs makes them ideal candidates for biosensors<sup>5–8</sup> and for bioimaging, drug delivery, and anticancer therapy.<sup>9,10</sup> However, despite their practical importance, major questions concerning the actual mechanisms of CD photoluminescence are still open,<sup>1</sup> especially the question of whether the fluorescence properties can be explained by single fluorophores or by aggregates.<sup>11</sup>

PAHs of various sizes have been used as model systems for studies on the luminescent behavior of CDs mentioned above.<sup>1,12,13</sup> Recently, Fu et al.<sup>14</sup> performed spectroscopic studies of model systems based on three PAHs, anthracene, pyrene, and perylene, embedded in a poly(methyl methacrylate) matrix to reproduce the features observed in actual CDs. An exciton self-trapping originating from relaxed excimer structures has been considered to be responsible for the large observed Stokes shifts. Even though these investigations point clearly to the important role of trapped excitons, a direct structural verification would be advantageous. The purpose of this work is to show, in combination with recent theoretical investigations of PAH dimers performed by our group<sup>15</sup> that these trapping effects induced by excimer formation indeed lead to significant Stokes shifts.

Inspired by the experimental observations of excitation-dependent photoluminescence, theoretical calculations, almost

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